

# Effect of $\text{Sm}_2\text{O}_3$ on the microstructure and electrical properties of $\text{SnO}_2$ -based varistors

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## Abstract

The effect of  $\text{Sm}_2\text{O}_3$  on the microstructure and non-linear electrical properties of (Co, Nb)-doped  $\text{SnO}_2$ -based varistors was investigated. The addition of  $\text{Sm}_2\text{O}_3$  improved the non-linear characteristics of (Co, Nb)-doped  $\text{SnO}_2$ -based varistors. The threshold electric field ( $E_B$ ) of  $\text{SnO}_2$ -based varistors increased significantly from 5340 to 12,460  $\text{V cm}^{-1}$  and the mean grain size decreased from 4.7  $\mu\text{m}$  to 1.7  $\mu\text{m}$  as  $\text{Sm}_2\text{O}_3$  concentration increased up to 0.20 mol%. There was an optimal value (0.20 mol%) of the  $\text{Sm}_2\text{O}_3$  concentration. The sample doped with 0.20 mol%  $\text{Sm}_2\text{O}_3$  had the highest non-linear coefficient ( $\alpha = 28$ ). The addition of  $\text{Sm}_2\text{O}_3$  reduced the relative density of (Co, Nb)-doped  $\text{SnO}_2$ -based varistors.

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**Keywords:** Varistors; Non-linear electrical behavior; Samarium oxide; Tin oxide; Rare earth; Ceramic

## 1. Introduction

Tin oxide is an n-type semiconductor with rutile crystalline structure and has many interesting electronic properties. In the recent decades, many papers have been published on new  $\text{SnO}_2$ -based varistors [1–12]. Tin oxide has very low densification rate due to the dominance of non-densifying mechanisms for mass transport such as surface diffusion or evaporation-condensation. Dense  $\text{SnO}_2$ -based ceramics can be achieved by introducing dopants (such as CoO,  $\text{MnO}_2$ ) [13] or by hot isostatic processing (HIP). The addition of cobalt oxide creates oxygen vacancies,  $\text{Co}'_{\text{Sn}}$  and  $\text{Co}''_{\text{Sn}}$  defects, which can segregate at grain boundary. Both defects of  $\text{Co}'_{\text{Sn}}$  and  $\text{Co}''_{\text{Sn}}$  can help the formation of Schottky barriers [1–12]. Pianaro et al. [14] found  $\text{SnO}_2$ -based ceramics doped with 1.00 mol% CoO and 0.05 mol%  $\text{Nb}_2\text{O}_5$  as promising varistor materials ( $\alpha \sim 8$ ,  $E_B = 1870 \text{ V cm}^{-1}$ ). A significant increase in  $\alpha$  of  $\sim 41$  and  $E_B = 4000 \text{ V cm}^{-1}$  is possible to be achieved by addition of 0.05 mol%  $\text{Cr}_2\text{O}_3$  [1]. Addition of  $\text{Cr}_2\text{O}_3$  increases the potential barrier values and the density of states at the grain boundary [1]. It was found that most of rare earth oxides such as  $\text{Gd}_2\text{O}_3$  [15],

$\text{Yb}_2\text{O}_3$  [16],  $\text{CeO}_2$  [17],  $\text{Pr}_2\text{O}_3$  [17,18],  $\text{Dy}_2\text{O}_3$  [19],  $\text{Er}_2\text{O}_3$  [20] and  $\text{La}_2\text{O}_3$  [3,4,17,21–23] could significantly improve the non-linear behavior of  $\text{SnO}_2$ -based varistors similar to that of  $\text{Cr}_2\text{O}_3$  [1,14]. Oliveira et al. [17] investigated the influence of  $\text{Pr}_2\text{O}_3$ ,  $\text{La}_2\text{O}_3$  and  $\text{CeO}_2$  on the structure and electrical properties of (Co, Nb, Cr)-doped  $\text{SnO}_2$ -based varistors. They reported that  $\text{Pr}_2\text{O}_3$  and  $\text{La}_2\text{O}_3$  segregate at the grain boundary and cause a considerable increase of  $\alpha$  value (to 62 and 81, respectively). The breakdown voltages of these varistors were 6866  $\text{V cm}^{-1}$  and 11,849  $\text{V cm}^{-1}$  for  $\text{Pr}_2\text{O}_3$  and  $\text{La}_2\text{O}_3$ , respectively. They reported that  $\text{CeO}_2$  as a dopant did not strongly influence the non-linear values of the  $\text{SnO}_2$ -based varistors ( $\alpha \sim 56$  and  $E_B = 5975 \text{ V cm}^{-1}$ ) as did  $\text{La}_2\text{O}_3$  and  $\text{Pr}_2\text{O}_3$ , indicating that  $\text{CeO}_2$  could have formed a solid state solution with  $\text{SnO}_2$  in the grains, but did not segregate sufficiently at the grain boundary [17]. Also, Oliveira et al. [21] investigated (Co, Nb, La)-doped  $\text{SnO}_2$  based varistors and found the values of  $\alpha$  and  $E_B$  as 46 and 12,432  $\text{V cm}^{-1}$ , respectively. Wang et al. investigated the (Pr, Co, Nb)-doped  $\text{SnO}_2$  varistors, and reported the non-linear coefficient of 61 and the electric field of 15,400  $\text{V cm}^{-1}$  for the varistors doped with 0.15 mol%  $\text{Pr}_2\text{O}_3$  [18]. Also, Wang et al. [15] reported  $\alpha \sim 24$  and  $E_B = 11,250 \text{ V cm}^{-1}$  for (Co, Nb, Gd)-doped  $\text{SnO}_2$ -based varistors. They further investigated the effect of  $\text{La}_2\text{O}_3$  on the non-linear behavior of (Co, Ta)-doped  $\text{SnO}_2$  and obtained  $\alpha \sim 23$  and  $E_B = 6350 \text{ V cm}^{-1}$  [23]. The

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values of  $\alpha \sim 26.3$  and  $E_B = 7422 \text{ V cm}^{-1}$  for (Co, Nb, Dy)-doped  $\text{SnO}_2$ -based varistors were obtained by Wang et al. [19]. Silva et al. [3] reported  $\alpha \sim 20$  and  $E_B = 4400 \text{ V cm}^{-1}$  for (Zn, Co, Ta)-doped  $\text{SnO}_2$  ceramics.

In this paper, the effect of  $\text{Sm}_2\text{O}_3$  on the microstructure and electrical properties of (Co, Nb)-doped  $\text{SnO}_2$  varistors was investigated.

## 2. Experimental procedures

The varistors were prepared by the standard ceramic technique. Commercial nano-sized  $\text{SnO}_2$  powders in the range of 50–80 nm were used (Fig. 1). The molar composition was  $(98.95 - x) \% \text{SnO}_2 + 1.00 \text{ CoO} + 0.05 \text{ Nb}_2\text{O}_5 + x \text{ Sm}_2\text{O}_3$  with  $x = 0.00, 0.05, 0.10, 0.20, 0.50$ , and  $1.00 \text{ mol\%}$ . The oxides used were analytical grade  $\text{SnO}_2$  (Merck),  $\text{CoO}$  (Riedel),  $\text{Nb}_2\text{O}_5$  (Merck) and  $\text{Sm}_2\text{O}_3$  (Alfa Aesar). The chemicals were wet milled in ethanol using zirconia balls for 1 hr. After drying at  $60^\circ\text{C}$ , the obtained powders were granulated with PVA binder and pressed into the discs with 12 mm in diameter and 1.0 mm in thickness by uniaxial pressing (15 MPa) followed by cold isostatic pressing (CIP) (240 MPa). After burning out of the PVA binder at  $650^\circ\text{C}$ , the discs were sintered in air at  $1300^\circ\text{C}$  for 1 h, and then cooled down to room temperature ( $5^\circ\text{C/min}$ ). In order to minimize or avoid cobalt loss during the sintering and to ensure the desired composition, the pellets were covered with their own powder. Silver paste was applied on the surfaces of the sintered discs to form electrodes by firing at  $600^\circ\text{C}$  for 10 min.

The apparent density of the sintered samples was measured by the Archimedes method. The crystal structure of the samples was characterized by X-ray diffraction (XRD) (Cu-K $\alpha$

radiation, Philips X-pert). From the XRD reflections corresponding to the (1 1 0) and (1 0 1) planes, a qualitative analysis of cell distortion in  $\text{SnO}_2$  samples was made. The signals were located in the angle values of  $2\theta \cong 26.61^\circ$  and  $2\theta \cong 33.89^\circ$ . The details can be found in [24,25]. The microstructure of the sintered pellets was observed using a TESCAN microscope. The grain size was calculated by the Mendelson equation [26].  $I$ – $V$  characteristics were measured using an ac power supply (Haefely Technology, Model CS 200). The frequency of the ac power supply was 60 Hz. The electrical non-linear coefficient ( $\alpha$ ) was obtained from:

$$\alpha = \frac{\log(I_2/I_1)}{\log(V_2/V_1)} \quad (1)$$

where  $V_1$  and  $V_2$  are voltages at the currents  $I_1$  and  $I_2$ , respectively. The  $\alpha$  values were obtained from the curves  $E$ – $J$  for current densities chosen between 1 and  $10 \text{ mA cm}^{-2}$ . The threshold electric field ( $E_B$ ) was obtained at the current density of  $1 \text{ mA cm}^{-2}$ .

## 3. Results and discussion

The XRD pattern of the  $\text{SnO}_2$  varistor doped with 0.50 mol%  $\text{Sm}_2\text{O}_3$  showed that all of the diffraction peaks could perfectly be indexed into a rutile  $\text{SnO}_2$  structure (JCPDS card, No. 41-1445). No other phase besides  $\text{SnO}_2$  was observed. The amount of additives used was very low and other possible phases may not have been detected because of the detection limit of the XRD equipment. Fig. 2 shows the SEM micrographs of system doped with 0.05 mol%  $\text{Sm}_2\text{O}_3$ . As clearly shown, the microstructure of this system is very simple and contains only one phase ( $\text{SnO}_2$  grains). Therefore, no

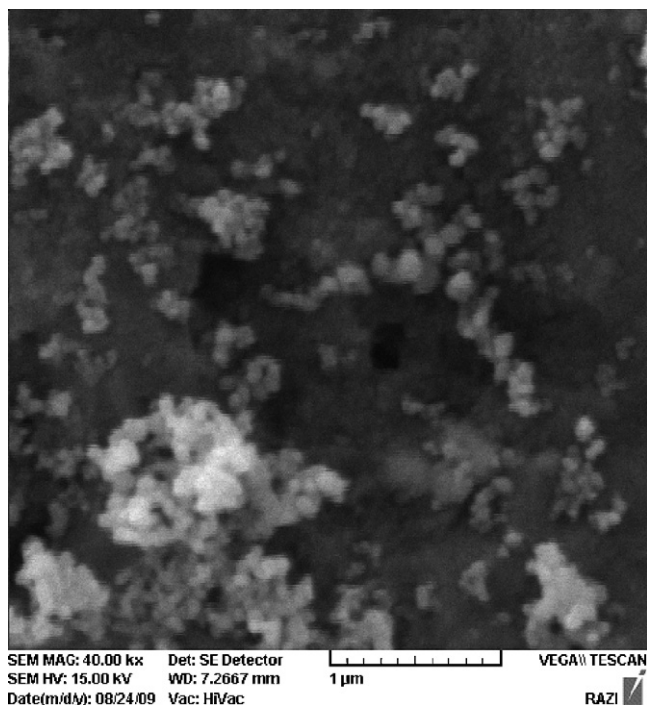


Fig. 1. SEM micrograph of the commercially nano-sized  $\text{SnO}_2$  powders.

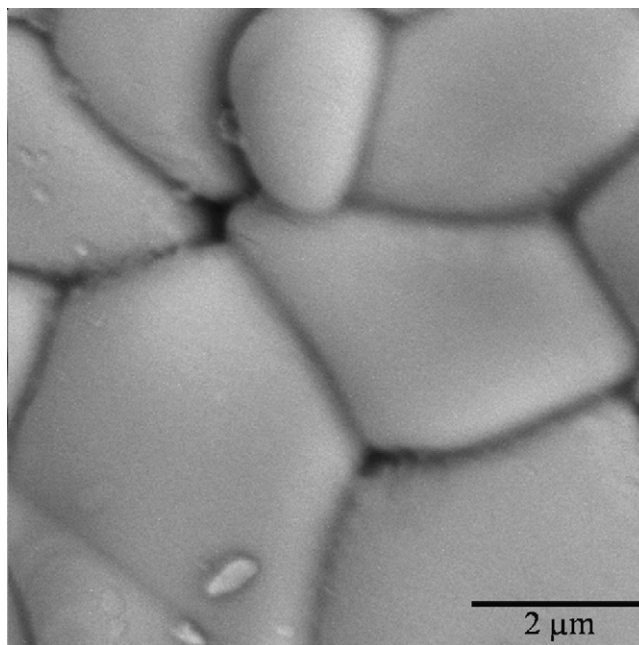


Fig. 2. SEM micrographs of (Co, Nb)-doped  $\text{SnO}_2$  based varistor doped with 0.20 mol%  $\text{Sm}_2\text{O}_3$  sintered at  $1300^\circ\text{C}$  for 1 h.

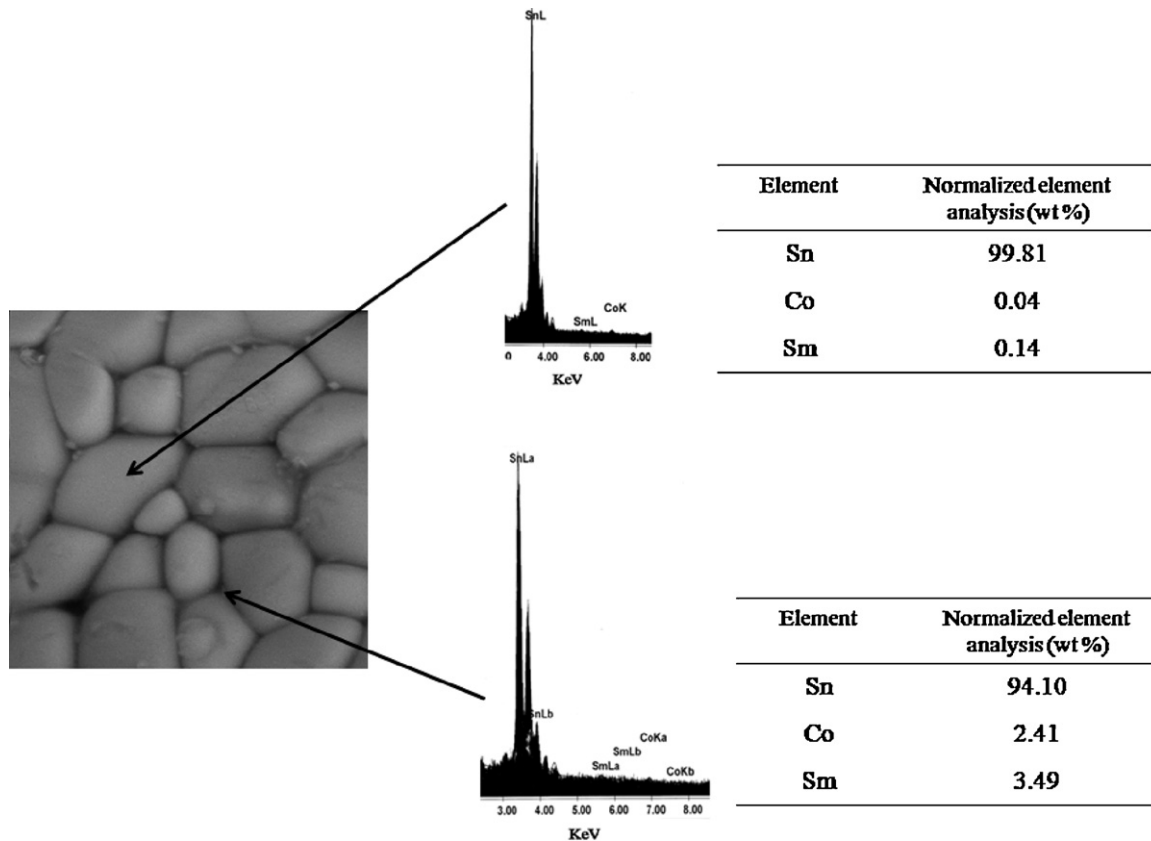


Fig. 3. Energy dispersive spectra (EDS) analysis for the sample doped with 0.50  $\text{Sm}_2\text{O}_3$  (mol%): grain (point 1); triple point (point 2).

secondary phase precipitated at the grain boundary was observed by XRD or SEM observation. This is in agreement with other studies [1–12,14–24].

Fig. 3 illustrates the energy dispersive spectra (EDS) analyses for the sample doped with 0.50  $\text{Sm}_2\text{O}_3$  (mol%) in different regions: grain (point 1), and triple point (point 2). As shown, in the point 1 of Fig. 3, the grains are predominantly composed of  $\text{SnO}_2$ . The weight percent ratios of Sn, Co and Sm are given in Fig. 3. Little amounts of samarium and cobalt were detected in the grain (Fig. 3, point 1) that could be due to the formation of solid solution in the doped  $\text{SnO}_2$ -based varistors. The amount of Co and Sm in the point 2 (triple point) increased as compared with the point 1 (grain) that could be attributed to the segregation of cobalt and samarium at this region. The presence of cobalt at the triple point may be due to the formation of a  $\text{Co}_2\text{SnO}_4$  precipitated phase, Similar to the findings of others [27]. The presence of samarium at the triple point (Fig. 3, point 2) from the EDS analysis could be attributed

to the segregation of samarium at the triple point and/or to the formation of some precipitates of  $\text{Sm}_2\text{Sn}_2\text{O}_7$ .

The characteristics of  $\text{SnO}_2$  varistor doped with different molar concentrations of  $\text{Sm}_2\text{O}_3$  are given in Table 1. The mean grain size of the  $\text{SnO}_2$ -based varistor without  $\text{Sm}_2\text{O}_3$  doping was  $4.70 \mu\text{m}$ . This value reduced to  $1.08 \mu\text{m}$  for the sample doped with 1.00 mol%  $\text{Sm}_2\text{O}_3$ . So, the addition of  $\text{Sm}_2\text{O}_3$  reduced the mean grain size of (Co, Nb)-doped  $\text{SnO}_2$  significantly and prevented the grains from growing, which is in agreement with the finding of other researchers [1–3,9,10,14–21].

The inhibition of grain growth is believed to occur by a mechanism of solute drag similar to the  $\text{CeO}_2$ -based ceramics doped with  $\text{Sm}_2\text{O}_3$ . The major driving forces leading to the segregation of an equilibrium concentration of solutes at the grain boundary are: (1) reduction in the elastic strain energy of the crystal lattice due to size difference between the solute and the host atoms for which it substitutes, and (2) the electrostatic

Table 1  
Characteristics of (Sm, Co, Nb)-doped  $\text{SnO}_2$ -based varistors.

$\text{Sm}_2\text{O}_3$ (mol%)	0.00	0.05	0.10	0.20	0.50	1.00
Density ( $\text{g cm}^{-3}$ )	6.87	6.84	6.80	6.78	6.75	6.73
Relative Density (%) <sup>a</sup>	98.85	98.42	97.84	97.55	97.12	96.83
Mean grain size ( $\mu\text{m}$ )	4.70	2.32	1.96	1.420	1.24	1.08
$\alpha$	10	17	21	28	16	13
$E_B$ ( $\text{V cm}^{-1}$ )	5340	6750	8470	12,460	10,400	9170

<sup>a</sup>Theoretical density of  $\text{SnO}_2$  is  $6.95 \text{ g cm}^{-3}$ .

potential of interaction between the aliovalent solutes and the charged grain boundary [28–30].

According to Table 1, the addition of Sm<sub>2</sub>O<sub>3</sub> reduced the relative density of (Co, Nb)-doped SnO<sub>2</sub>-based varistors. This value was 98.85% of the theoretical density (TD) for (Co, Nb)-doped SnO<sub>2</sub> system. The addition of Sm<sub>2</sub>O<sub>3</sub> up to 1.00 mol% reduced the relative density to 96.83% TD (Table 1). The increase in porosity is probably due to the segregation of Sm<sub>2</sub>O<sub>3</sub> at the grain boundary, hindering the sintering process. Sm<sub>2</sub>O<sub>3</sub> appeared to exert an influence on the microstructure of (Co, Nb)-doped SnO<sub>2</sub>-based varistors similar to that observed by Pianaro et al. [1] in (Co, Nb, Cr)-doped SnO<sub>2</sub>-based varistors.

The threshold electric field versus the current density of the SnO<sub>2</sub>-based varistors doped with Sm<sub>2</sub>O<sub>3</sub> is illustrated in Fig. 4. As shown, the threshold electric field of the SnO<sub>2</sub>-based varistors increased from 5340 to 12,460 V cm<sup>−1</sup> as Sm<sub>2</sub>O<sub>3</sub> concentration increased to 0.20 mol%. The highest non-linear coefficient ( $\alpha = 28$ ) was obtained for the sample doped with 0.20 mol% Sm<sub>2</sub>O<sub>3</sub>. The improvement of non-linear characteristics of (Co, Nb, Sm)-doped SnO<sub>2</sub>-based varistors can be due to the segregation of Sm<sub>2</sub>O<sub>3</sub> at the grain boundary, probably inducing electronic interface states that can trap states at the SnO<sub>2</sub>–SnO<sub>2</sub> interface. Similar results were found by other researchers in SnO<sub>2</sub>-based varistors doped with other transition metal oxides [1–24]. Sm<sub>2</sub>O<sub>3</sub> apparently is present at the grain boundary and may alter the potential barrier at the grain boundary, affecting the varistor properties.

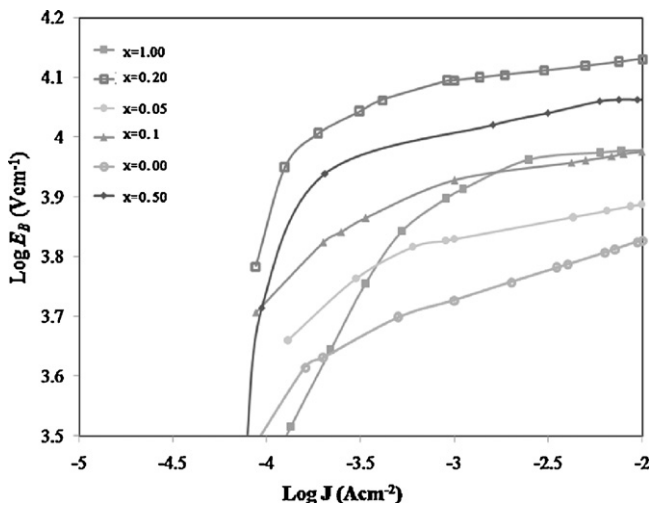


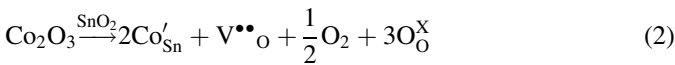
Fig. 4. Applied electric field versus current density for (Sm, Co, Nb)-doped SnO<sub>2</sub> varistors doped with different concentrations of Sm<sub>2</sub>O<sub>3</sub>.

Table 2  
Lattice parameters of the sintered samples.

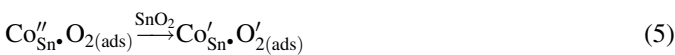
System	Composition (mol%)	<i>a</i> ± 0.00001 (Å)	<i>c</i> ± 0.00001 (Å)	<i>V</i> = <i>a</i> <sup>2</sup> <i>c</i> ± 0.00001 (Å <sup>3</sup> )
S	100.00 SnO <sub>2</sub>	4.73088	3.18186	71.21406
SC	99.00 SnO <sub>2</sub> + 1.00 CoO	4.73088	3.18485	71.28106
SCN	98.95 SnO <sub>2</sub> + 1.00 CoO + 0.05 Nb <sub>2</sub> O <sub>5</sub>	4.72913	3.18504	71.23238
SCNSm	98.75 SnO <sub>2</sub> + 1.00 CoO + 0.05 Nb <sub>2</sub> O <sub>5</sub> + 0.20 Sm <sub>2</sub> O <sub>3</sub>	4.73063	3.18781	71.33947

Addition of Sm<sub>2</sub>O<sub>3</sub> in the concentration exceeding 0.20 mol% degraded the non-linear properties of SnO<sub>2</sub> varistors mainly due to the fact that the microstructure of (Co, Nb, Sm)-doped SnO<sub>2</sub>-based varistors depended on the concentration of Sm<sub>2</sub>O<sub>3</sub>. The large amounts of Sm<sub>2</sub>O<sub>3</sub> might decrease the grain boundary mobility leading to a decrease in density and consequently, increasing the porosity of (Co, Nb, Sm)-doped SnO<sub>2</sub>-based varistors. Also, the higher concentrations of samarium at triple point measured by EDS analysis (Fig. 3, point 2) could be related to the segregation of samarium and/or the formation of Sm<sub>2</sub>Sn<sub>2</sub>O<sub>7</sub> at the triple point. When precipitates are present in higher concentrations at triple points, as in the present case, due to the addition of 0.5 mol% Sm<sub>2</sub>O<sub>3</sub>, they can affect the nature of SnO<sub>2</sub>–SnO<sub>2</sub> junctions and, hence, the number of active barriers. Simões et al. [12] studied the effect of excess precipitates on the non-linear behavior of (Co, Ta, La)-doped SnO<sub>2</sub>-based varistors. They proposed that an excessive amount of precipitates may be deleterious to the non-ohmic properties of SnO<sub>2</sub>-based varistors because they may create adjacent regions at grain boundary with lower concentrations of segregated metal atoms.

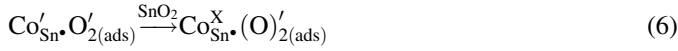
The lattice parameters of undoped and doped SnO<sub>2</sub> with CoO, Nb<sub>2</sub>O<sub>5</sub> and Sm<sub>2</sub>O<sub>3</sub> determined from the XRD patterns are listed in Table 2. Addition of cobalt increased the lattice volume of SC system in comparison with undoped SnO<sub>2</sub> (Table 2, S system). This could be due to the formation of solid solution. This result is in accordance with the findings of other studies [22,31,32]. Also, the high densification observed in the CoO-doped SnO<sub>2</sub> can be explained by the substitution of Sn<sup>4+</sup> (0.071 nm) by Co<sup>2+</sup> (0.078 nm) or by Co<sup>3+</sup> (0.063 nm) in the SnO<sub>2</sub> crystalline lattice according to:



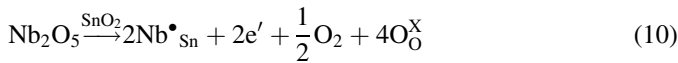
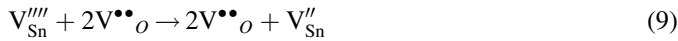
According to the above equations, cobalt could stabilize as Co<sup>2+</sup> or Co<sup>3+</sup> [1,14,31]. The increase in the lattice volume of un-doped SnO<sub>2</sub> (Table 2, SC system) could be attributed to the stabilizing of cobalt ions as Co<sup>2+</sup>. The cobalt oxide could influence the adsorption of oxygen species at the grain boundary as represented by the Eqs. (4)–(7) [33]:





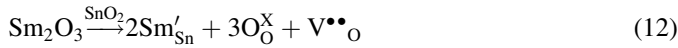


A decrease in the lattice volume of Co-doped  $\text{SnO}_2$  was observed by the addition of  $\text{Nb}_2\text{O}_5$  (Table 2, SCN system) could be related to the formation of solid solution. This is in agreement with the findings of other researchers [22,31]. The substitution of  $\text{Sn}^{4+}$  by  $\text{Nb}^{5+}$  can be written as the following reactions [1,14,22]:



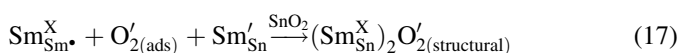
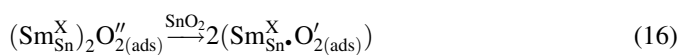
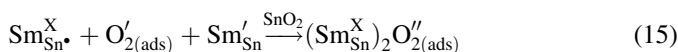
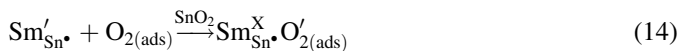
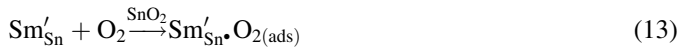
$\text{Nb}_2\text{O}_5$  is an electron donor that increases the electrical conductivity of  $\text{SnO}_2$ . The electrons introduced through Eq. (10) may be involved in the transformation of  $\text{V}^{\bullet\bullet}_{\text{O}}$  into  $\text{V}^{\bullet}_{\text{O}}$  [34].

The increase in both lattice parameters,  $a$  and  $c$ , in the sample doped with 0.20 mol%  $\text{Sm}_2\text{O}_3$  (Table 2, SCNSm system) could be attributed to the substitution of  $\text{Sn}^{4+}$  by  $\text{Sm}^{3+}$ . Since the ionic radius of  $\text{Sm}^{3+}$  (0.096 nm) is larger than that of  $\text{Sn}^{4+}$  (0.071 nm), the substitution of  $\text{Sn}^{4+}$  by  $\text{Sm}^{3+}$  as [35]:



leads to the distortion of  $\text{SnO}_2$  lattice. Due to the partial substitution of  $\text{Sn}^{4+}$  by  $\text{Sm}^{3+}$ , much more  $\text{Sm}_2\text{O}_3$  will reside at the  $\text{SnO}_2$  grain boundary. The species of  $\text{Sm}'_{\text{Sn}}$  segregated at the  $\text{SnO}_2$  grain boundary are the origin of the potential barrier. They create oxygen vacancies to promote adsorbed oxygen species formation, which are responsible for the barrier formation near the grain boundary.

Similar to  $\text{Cr}_2\text{O}_3$ ,  $\text{La}_2\text{O}_3$  and other rare earth oxides [3,9,10,23], samarium oxide could react with oxygen species according to the following equations:



The role of  $\text{Co}''_{\text{Sn}}$ ,  $\text{Sm}'_{\text{Sn}}$ ,  $\text{Sn}''_{\text{Sn}}$ ,  $\text{V}''''_{\text{Sn}}$  and  $\text{V}''_{\text{Sn}}$  defects (due to the reduction of  $\text{Sn}^{4+}$  to  $\text{Sn}^{2+}$  in the interstitial sites) would be related to the increase of the  $\text{O}''_2$ ,  $\text{O}'$  and  $\text{O}''$  adsorbed at the grain boundary interface, causing a decrease in the conductivity of grain boundary by the donation of electrons to the  $\text{O}_2$  adsorbed at the grain boundary surface. As pointed out in Ref. [33], the

adsorbed oxygen at the grain boundary captures electron from the negatively charged defects at the grain boundary and stays at the interface [33]. In polycrystalline semiconductors, the trapping of electron at the grain boundary has a decisive influence on the electrical transport properties by means of the formation of electrostatic potential barriers [36,37].

Recently, Bueno et al. [38] published a review on the voltage-dependent resistance feature. They discussed completely the details of the basic physical principles involved with the non-ohmic properties of both  $\text{SnO}_2$  and  $\text{ZnO}$ -based non-ohmic systems, with an evaluation of contribution of the dopants to the electronic properties, to the final microstructure and consequently to non-ohmic behavior of the systems. The main conclusion reached by the authors (in the mentioned review) was that despite the significant difference in the microstructure nature of  $\text{ZnO}$  and  $\text{SnO}_2$  varistors, their physical nature is the same that can be described by a Schottky-type double barrier, which is presumed to be associated to oxygen species at the grain boundary of both  $\text{ZnO}$  and  $\text{SnO}_2$ -based varistors [36–39]. Glot et al. studied the non-ohmic conduction in  $\text{SnO}_2$ -based varistors [40,41]. They found high values of the activation energy of electrical conduction and the barrier height in  $\text{SnO}_2$ -based varistors and concluded the conduction process to be thermally activated not only at low electric field but also at high electric field. This makes possible a simple explanation of the conduction in  $\text{SnO}_2$ -based varistors as thermoionic emission across a barrier height dependent on electric field [41].

#### 4. Conclusions

The microstructure and electrical properties of (Co, Nb, Sm)-doped  $\text{SnO}_2$  based varistors were evaluated. The non-linear properties of (Co, Nb)-doped  $\text{SnO}_2$ -based varistors increased by the addition of  $\text{Sm}_2\text{O}_3$ . The optimum  $\text{Sm}_2\text{O}_3$  concentration of (Sm, Co, Nb)-doped  $\text{SnO}_2$ -based varistors was determined as 0.20 mol%. The highest non-linear coefficient ( $\alpha = 28$ ) was obtained for the sample doped with 0.20 mol%  $\text{Sm}_2\text{O}_3$ . The results indicated that  $\text{Sm}_2\text{O}_3$  may segregate at the grain boundary, and induce electronic interface states that can trap charge at the  $\text{SnO}_2$ – $\text{SnO}_2$  interface and lead to the increase of the  $\alpha$  value. The threshold field of this sample was  $12,460 \text{ V cm}^{-1}$ .

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