

Short communication

Size–strain line broadening analysis of high-permittivity oxynitrides
BaTaO₂N and BaNbO₂N

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Abstract

X-ray line broadening analysis of polycrystalline BaTaO₂N and BaNbO₂N was performed to extract microstructural information beyond the average structural descriptions. Both BaTaO₂N and BaNbO₂N display markedly broader diffraction peak profiles than observed from an oxide analogue KTaO₃. Examination of the angular dependence of the integral peak breadth shows that the line broadening of BaTaO₂N is caused by both finite crystallite size ($D_v = 117$ nm) and micro-strain ($\varepsilon = 0.038\%$), whereas the latter effect dominates in the case of BaNbO₂N ($D_v = 574$ nm, $\varepsilon = 0.11\%$). The above results indicate that the temperature-independent high-permittivity behavior of BaTaO₂N arises due to local atomic displacements, and further imply that even stronger dielectric polarizations could be generated in BaNbO₂N.

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1. Introduction

Being comparable in size, charge, and electronegativity, the two anions O^{2−} and N^{3−} have similar bonding preferences and can form oxynitride type mixed anion compounds. As the coexistence of O^{2−} and N^{3−} in the anion sub-lattice gives rise to physical and chemical properties clearly distinct from those of either the pure oxide or the pure nitride relatives, the oxynitrides are regarded a promising system for exploring unconventional material performances. In fact, there have been considerable research efforts in the past decade to develop novel functional materials using various transition metal oxynitrides [1–4]. As an example, the perovskite BaTaO₂N was reported to exhibit unusually high and temperature-independent dielectric constants (κ) both in polycrystalline ($\kappa \approx 5000$ at 180–300 K) and thin film ($\kappa \approx 220$ at 4–300 K) forms [3,5]. Interestingly BaTaO₂N has an average crystal structure of primitive cubic symmetry, which is seemingly inconsistent with high- κ behavior. Later studies, however, revealed that the BaTaO₂N lattice contains substantial local distortions and is largely deviated from the ideal cubic geometry [6–8]. Such a local structural phenomenon is believed to derive from the random O/N distribution in the anion sub-lattice, coupled with the inequity of Ta–O and Ta–N bonds.

In order for better utilizing the dielectric property of BaTaO₂N, it is of crucial importance to comprehend the nature of the lattice polarization. In this regard, we examined the microstructural aspects of BaTaO₂N, in comparison with other cubic perovskites BaNbO₂N and KTaO₃, by powder diffraction line profile analysis (LPA) based on the simplified integral breadth method. Since the earlier recognition by Scherrer of the dependence of diffraction line width on crystallite size [9], diffraction LPA has been widely used for studying microstructural characteristics, *i.e.* size and strain, of various extended solids [10]. While there are known several different LPA schemes in practice, integral breadth methods are frequently used for quick assessment of the size–strain line broadening effects that are caused by the finite size of the coherently scattering domains and by lattice strain [11]. As is clearly found from the present LPA and electron microscopy, BaTaO₂N and BaNbO₂N are formed from much smaller crystallites and with larger micro-strains than the oxide KTaO₃. Moreover it is revealed that BaNbO₂N contains larger crystallites and larger strain than BaTaO₂N.

2. Experimental

Powder samples of BaTaO₂N and BaNbO₂N were prepared by ammonolysis reactions using BaCO₃ (Aldrich, 99.98%), Ta₂O₅ (Aldrich, 99.99%), and Nb₂O₅ (Aldrich, 99.99%). For each composition, a stoichiometric mixture of powder reagents

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was heated at 1273 K for a total of 60 h with intermittent grinding, in an ammonia flow of ≈ 50 sccm. A reference oxide KTaO_3 was also prepared from a stoichiometric mixture of K_2CO_3 (Alfa-Aesar, 99.0%) and Ta_2O_5 (Aldrich, 99.99%) by heat treatment in air at 1273 K for 60 h.

Synchrotron powder X-ray diffraction (XRD) data were collected at 295 K, on a high-resolution powder diffractometer (BL-8C2) installed at Pohang Accelerator Laboratory (PAL) in Korea. The use of high-resolution synchrotron radiation ($\lambda = 1.550 \text{ \AA}$) helped to minimize any instrumental contributions to the line broadening effect. A standard reference material, LaB_6 (NIST SRM 660a, $a = 4.156468 \text{ \AA}$) [12], was employed for wavelength calibration and also for correcting instrumental broadening of line widths. Diffraction profiles were recorded over the 2θ range $10\text{--}130^\circ$ with a step size of 0.01° for BaTaO_2N , BaNbO_2N , and KTaO_3 , and 0.002° for LaB_6 . The sample holder was spun at 60 rpm in the Bragg–Brentano configuration.

The grain sizes and morphologies of the sample powders were examined using field-emission scanning electron microscopes (Hitachi, S-4800 and S-4200). SEM powder samples were mounted on aluminum stubs using double-sided conductive carbon tabs, and coated with Pt prior to imaging.

3. Results and discussion

Synchrotron XRD profiles of BaTaO_2N , BaNbO_2N , KTaO_3 , and LaB_6 are presented in Fig. 1. The four patterns commonly correspond to the simple cubic lattice, without any indication of peak splitting or any impurity phases. Rietveld refinement confirmed that all the samples belong to the $Pm\bar{3}m$ space group with lattice parameters of $4.11225(2) \text{ \AA}$, $4.12131(3) \text{ \AA}$, and $3.98849(3) \text{ \AA}$ for BaTaO_2N , BaNbO_2N , and KTaO_3 , respectively. Each sample presented 14–22 diffraction peaks at $2\theta < 130^\circ$ ($d > 0.86 \text{ \AA}$), for which the full width at half maximum (Γ , Fig. 2a) and the Lorentz fraction (η) were determined by fitting a pseudo-Voigt function. Subsequently the integral breadth (β) was calculated using the equation [13]:

$$\beta = \frac{\Gamma}{2} \left\{ \sqrt{\frac{\pi}{\ln 2}} + \eta \left(\pi - \sqrt{\frac{\pi}{\ln 2}} \right) \right\}$$

After correcting for the instrumental broadening, β_{inst} , with LaB_6 as a reference, Williamson–Hall (W–H) plots [14] for KTaO_3 , BaTaO_2N , and BaNbO_2N were constructed based on the function:

$$(\beta_i - \beta_{\text{inst}}) \cos \theta_i = \frac{\lambda}{D_v} + 4\epsilon \sin \theta_i$$

where β_i is the integral breadth (in radians 2θ) of the i th Bragg reflection positioned at $2\theta_i$ and λ is the radiation wavelength. The slope and ordinate intercept of the W–H plot (Fig. 2b) were used to quantify crystallite size (D_v) and strain (ϵ) contributions to the peak broadening.

Although the three perovskite samples were prepared using the same sintering temperature and time, the resulting microstructural characteristics turned out to be quite distinct. While KTaO_3 consists of crystallites in the micron range with

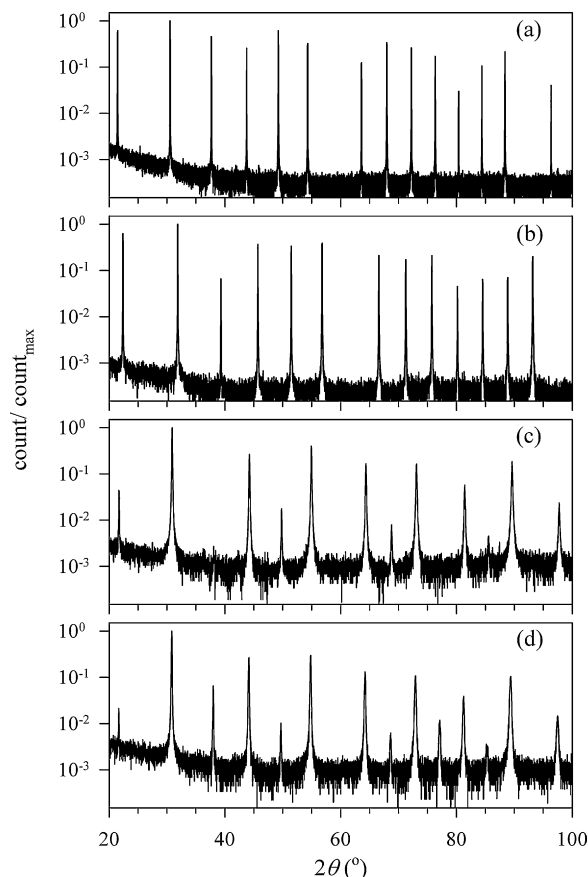


Fig. 1. Synchrotron XRD ($\lambda = 1.550 \text{ \AA}$) patterns for (a) LaB_6 , (b) KTaO_3 , (c) BaTaO_2N , and (d) BaNbO_2N .

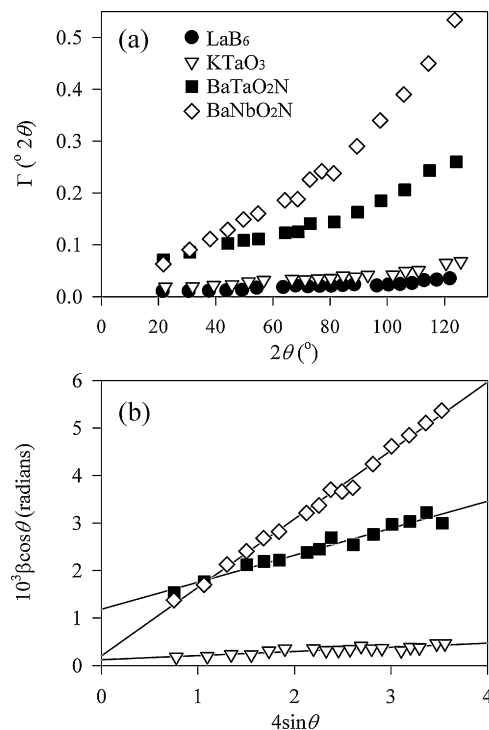


Fig. 2. (a) Peak widths (Γ) for LaB_6 , KTaO_3 , BaTaO_2N , and BaNbO_2N , and (b) the corresponding Williamson–Hall plots. The same legend symbols are used for both panels.

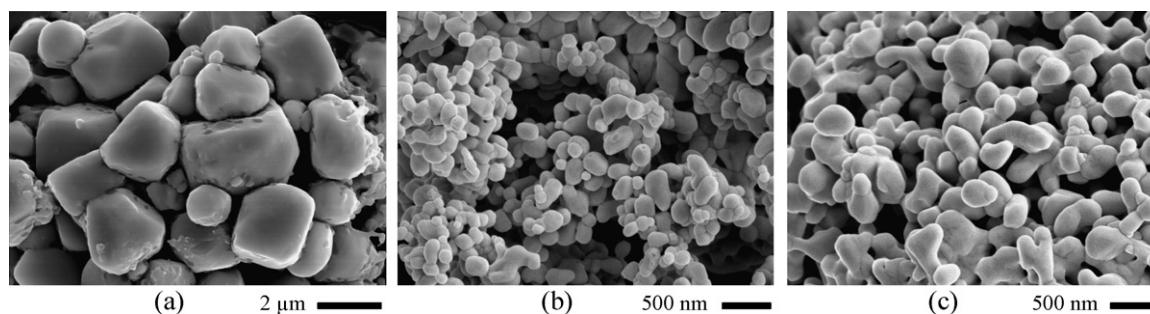


Fig. 3. SEM images of (a) KTaO_3 , (b) BaTaO_2N , and (c) BaNbO_2N .

minimal strain ($D_v = 1034$ nm, $\varepsilon = 0.005\%$), both oxynitride samples have markedly smaller crystallite sizes and larger strains (BaTaO_2N : $D_v = 117$ nm, $\varepsilon = 0.038\%$; BaNbO_2N : $D_v = 574$ nm, $\varepsilon = 0.11\%$). The crystallite size and strain are directly affected by the degree of crystal imperfection. That is, the higher concentration of lattice defects should reduce the crystallite size while increasing the strain. The present LPA shows that the two oxynitride type samples are more defective than the oxide example KTaO_3 . The better sintering of the latter, compared with the two oxynitrides, may be partly due to the flux behavior of the alkali component, and also to the generic refractory nature of the nitride-containing phases. Additionally, it should be mentioned that the intrinsic local structural distortion in oxynitrides contributes to develop substantial amounts of micro-strain within the crystallites. Recent experimental investigations using extended X-ray absorption fine structure [6], pair-distribution-functions [7], and electron diffraction [8] have elucidated that BaTaO_2N consists of irregularly distorted TaO_4N_2 octahedra, which must have a profound impact on the lattice strain. It is highly probable that such a structural feature exists universally in the oxynitride perovskite system, and we regard the size–strain line broadening of oxynitride perovskites as being associated with compositional aspects and structural disorders as well.

An interesting finding is that both the crystallite size and strain are assessed as being significantly larger for BaNbO_2N than for BaTaO_2N . Taking into account that the niobium-related phases tend to melt at lower temperatures than their tantalum analogues (Ta: 3290 K; Nb: 2750 K; Ta_2O_5 : 2145 K; Nb_2O_5 : 1785 K), it is duly expected that BaNbO_2N would have a lower melting point than BaTaO_2N and would undergo faster crystallite growth. However the larger strain of BaNbO_2N is not explainable by the above reasoning, but is attributable to the more pronounced local structural disorder, as described below.

For BaTaO_2N , the packing incompatibility was pointed out as a plausible origin of the local atomic displacements [8]. The perovskite tolerance factor for the composition BaTaO_2N is calculated as 1.04, which practically means that the size of a Ta^{+5} ion is too small to fit the octahedral cavity. Accordingly, the bond valence sum (BVS) calculations [15] for the average structure of BaTaO_2N reveal a notable underbonding of the Ta ion. Using the software VALENCE [16], we obtain +2.71 valence units (v.u.) for Ba, +4.54 v.u. for Ta, −2.13 v.u. for O, and −2.99 v.u. for N ions, and subsequently a global instability index (GII) of 0.388 for BaTaO_2N . A GII

value exceeding 0.2 is generally considered unrealistic for any standard state structures [15], and this BVS calculation itself can be a strong implication of structural relaxation in BaTaO_2N . Previously it was rationalized that such an unfavorable bonding geometry of BaTaO_2N can be alleviated by the displacement of Ta ions away from the ideal high-symmetry site [8]. Similarly for the average structure of BaNbO_2N , the BVS calculation yields +2.67 v.u. for Ba, +4.41 v.u. for Nb, −2.10 v.u. for O, and −2.95 v.u. for N ions, and a GII of 0.402. Here the bond valence of the octahedral cation and the GII for BaNbO_2N deviate from their respective ideal values, even farther than in BaTaO_2N . This leads us to the deduction that BaNbO_2N would contain more prominent lattice disorders than BaTaO_2N . It is also noteworthy that Nb, being more electronegative than Ta, is more prone to intra-octahedral distortions *via* second-order Jahn–Teller distortions [17,18]. We regard LPA results as corroborating the above assumptions well.

The crystallite morphologies and sizes of KTaO_3 , BaTaO_2N , and BaNbO_2N are compared using SEM in Fig. 3. Although the crystallite sizes are somewhat distributed in all three cases, the majority of the observed crystallites fall in to size ranges roughly corresponding to the LPA result: 1–3 μm for KTaO_3 , 80–300 nm for BaTaO_2N , and 300–500 nm for BaNbO_2N . It is noted that the particulates of KTaO_3 maintain cube-like polyhedral shapes whereas those of BaTaO_2N and BaNbO_2N have rather curvy surfaces. The loss of crystal habits in the above oxynitride powders is indicative of the lack of long-range lattice periodicity, namely the prevalence of irregular structural distortions in local scales. Such a morphological characteristic is in accordance with the presence of large micro-strains. In view of the dielectric property, the larger strain of BaNbO_2N gives promises as a new high- κ material.

4. Conclusions

By X-ray line broadening analysis, it was demonstrated that the oxynitrides BaTaO_2N and BaNbO_2N are formed with smaller crystallite sizes and greater micro-strains than the oxide analogue KTaO_3 . It is judged that the micro-strains of the above oxynitrides stem largely from the local atomic displacements, which has immediate relevance to the dielectric behavior. Based on the present LPA results, the unusual dielectric properties of the simple cubic perovskite BaTaO_2N can be attributed to the local polarizations induced by geometry

relaxation. Furthermore, BaNbO₂N is identified as a promising high-permittivity candidate.

Acknowledgements

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