

# Magnetic behavior of nickel–bismuth ferrite synthesized by a combined sol–gel/thermal method

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## Abstract

The bulk  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$  ferrites with various  $\text{Bi}^{3+}$  concentration ( $x=0, 0.1, 0.15$ ) were synthesized via sol–gel procedure, starting from nickel, bismuth and iron nitrate powders, followed by the conventional thermal treatment. The structural and magnetic properties of the as-prepared ferrites were studied by means of X-ray diffraction, alternating gradient force magnetometry and Faraday balance. The anisotropy constant was determined by the law of approach to saturation (LAS) model. An increasing  $\text{Bi}^{3+}$  concentration in  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$  leads to a decrease in the saturation magnetization, Néel temperature and the anisotropy constant of the material. © 2012 Elsevier Ltd and Techna Group S.r.l. All rights reserved.

**Keywords:** Nickel–bismuth ferrite; Combined sol–gel/thermal synthesis; Magnetic behavior; Anisotropy constant

## 1. Introduction

Ferrites with the spinel structure of the type  $\text{MFe}_2\text{O}_4$  (M is a divalent cation) are magnetic ceramics of great importance [1]. Their applications range from simple function devices to sophisticated devices for electronics. Some interesting applications are in computer peripherals, telecommunication equipments, electronic and microwave devices, magnetic media, recording devices and magnetic cards. It is well recognized that electrical and magnetic properties of spinel ferrites depend on various parameters such as their processing conditions and chemical composition [1–4]. These complex oxides crystallize in an *fcc* structure with the transition metal cations ( $\text{M}^{2+}$ ,  $\text{Fe}^{3+}$ ) residing at the interstices formed by oxygen anions. To emphasize the site occupancy at the atomic level, the structural formula of spinel ferrites may be written as  $(\text{M}_1^{2+}{}_\lambda\text{Fe}_\lambda^{3+})[\text{M}_\lambda^{2+}\text{Fe}_2^{3+}]_2\text{O}_4$ , where parentheses and square brackets denote cation sites of tetrahedral (A) and octahedral [B] coordinations, respectively.  $\lambda$  represents

the degree of inversion defined as the fraction of the (A) sites occupied by  $\text{Fe}^{3+}$  cations.

Nickel ferrite,  $\text{NiFe}_2\text{O}_4$ , as a soft magnetic n-type semiconducting material, is an important member of the spinel family with a fully inverse structure ( $\lambda=1$ ) in the bulk state [5,6]. It has been revealed that an increase in temperature and the particle size reduction of a spinel induce the cation redistribution that is directed towards random arrangement [7]. In both, undoped and substituted forms,  $\text{NiFe}_2\text{O}_4$  serves as a material for high-frequency applications in telecommunication field. Due to its high resistivity and low eddy current losses it is used in radio-frequency circuits, high quality filters, rod antennas, transformer cores, read/write heads for high-speed digital tape and operating devices [8].

In previous paper [9], we reported structural and magnetic properties of  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$  nanoparticles prepared via sol–gel method. The present work represents a continuation of our work [9] with the goal to characterize the bulk  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$  ( $x=0, 0.1$ , and  $0.15$ ) synthesized by the combined sol–gel/thermal method. Crystal structure and magnetic properties (Néel temperature, saturation magnetization and anisotropy constant) of the as-prepared bulk samples are investigated as a function of  $\text{Bi}^{3+}$  concentration, and are

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<sup>1</sup>This work is dedicated to my wife Ghazale.

compared with those of  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$  nanoparticles. Note that although a lot of work concerning the substituted nickel ferrites has already been published, see, e.g. [4,9–18], to the best of our knowledge, there is no report in the literature on  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$  prepared via sol–gel procedure followed by the thermal treatment.

## 2. Experimental

Nickel–bismuth ferrites,  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$  with  $x=0, 0.1$ , and  $0.15$ , were synthesized in the bulk form by the combined sol–gel/thermal method. At first, stoichiometric amounts of  $\text{Fe}(\text{NO}_3)_3$ ,  $\text{Ni}(\text{NO}_3)_2$  and  $\text{Bi}(\text{NO}_3)_3$  were dissolved completely in deionized water. In the dissolving process, the molar ratio of cation concentrations of  $\text{Ni}^{2+}/(\text{Bi}^{3+} + \text{Fe}^{3+})$  was fixed at 1:2 and that of  $x\text{Bi}^{3+}/(2-x)\text{Fe}^{3+}$  was varied with  $x=0, 0.1$  and  $0.15$ . Each aqueous solution containing  $\text{Ni}^{2+}$ ,  $\text{Fe}^{3+}$  and  $\text{Bi}^{3+}$  was poured into citric acid with the ratio of  $(\text{Ni}^{2+} + \text{Bi}^{3+} + \text{Fe}^{3+})/(\text{citric acid})=3$ . The mixtures were stirred and slowly evaporated at 353 K to form gels. These gels were dried at 373 K for 2 h and then annealed in air at 1273 K for 3 h.

The structure of the as-prepared  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$  was studied using an X-ray diffractometer (D8, Bruker, Germany) with  $\text{CuK}\alpha$  radiation ( $\lambda_{\text{Cu}}=1.54056 \text{ \AA}$ ). The Néel temperature and the hysteresis ( $M-H$ ) loops of the samples were measured by a Faraday balance equipped with a permanent magnet (0.02 T) and an alternating gradient force magnetometer (AGFM), respectively.

## 3. Results and discussion

The XRD patterns of the as-prepared nickel–bismuth ferrites with various  $\text{Bi}^{3+}$  concentration ( $x=0, 0.1$ , and  $0.15$ ) are shown in Fig. 1. The diffraction patterns of the samples confirm that the spinel phase was formed. It is clearly seen that the diffraction lines for all three samples are relatively narrow, indicating their polycrystalline (bulk) character. The lattice parameters of the cubic ferrites,  $a$ , is calculated using the following formula [19]:

$$a = \frac{\lambda_{\text{Cu}} \sqrt{h^2 + k^2 + l^2}}{2 \sin \theta} \quad (1)$$

where  $\theta$  is the Bragg angle and  $h$ ,  $k$  and  $l$  are the Miller indices. Variation of the lattice parameter of the ferrites with  $\text{Bi}^{3+}$  concentration is presented in Fig. 2. As clearly visible, the lattice parameter of  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$  samples increases with increasing  $\text{Bi}^{3+}$  concentration. This can be attributed to the different sizes of  $\text{Bi}^{3+}$  and  $\text{Fe}^{3+}$  ions;  $\text{Bi}^{3+}$  has larger ionic radius (1.31 Å) than  $\text{Fe}^{3+}$  ions (0.66 Å) [20].

From the determined  $a$  values, X-ray density ( $d_x$ ) of the Bi-substituted nickel ferrites was calculated using the following formula [19]:

$$d_x = \frac{8M}{Na^3} \quad (2)$$

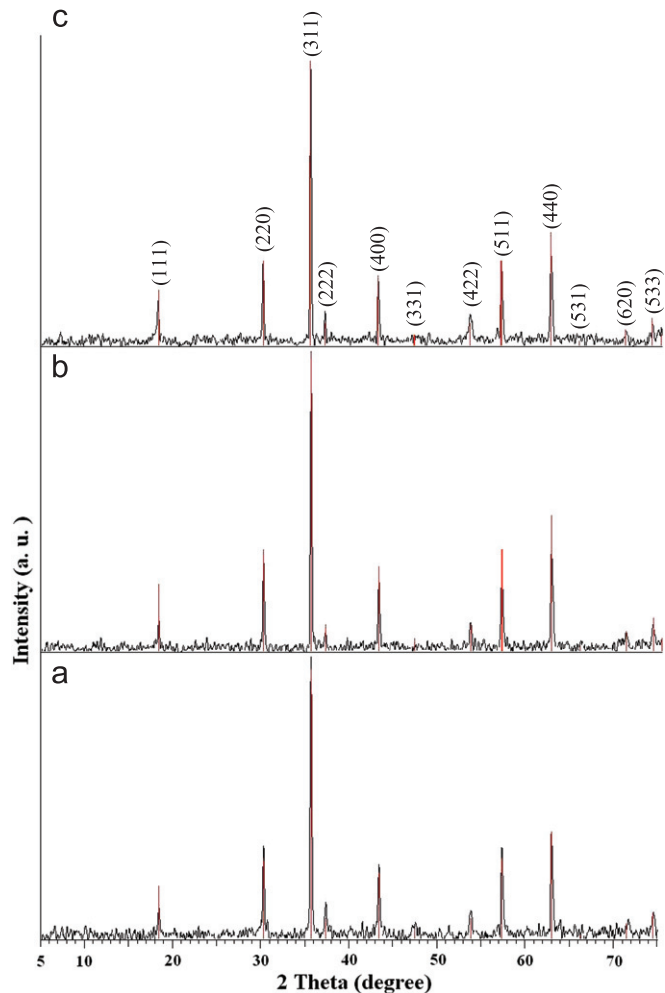


Fig. 1. XRD patterns of polycrystalline  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$  samples: (a)  $x=0$ , (b)  $x=0.1$ , (c)  $x=0.15$ , prepared by the thermal treatment of  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$  nanoparticles synthesized by the sol–gel procedure. The diffraction peaks are denoted by Miller indices corresponding to the spinel phase (JCPDS PDF 10-0325).

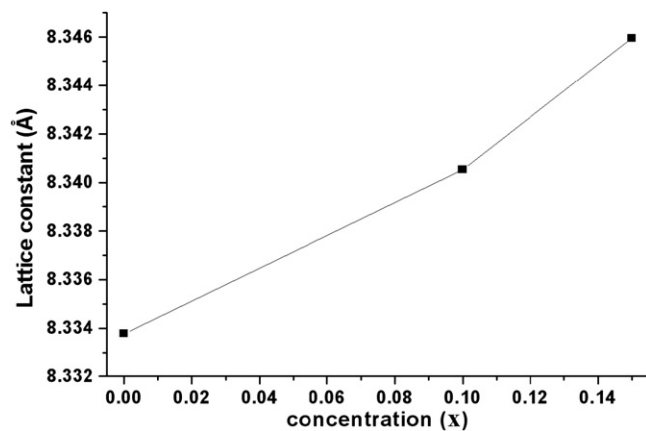


Fig. 2. Variation of the lattice parameter of  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$  with Bi content.

where  $M$  and  $N$  are the molecular weight and Avogadro's number, respectively. The  $d_x$  values are presented in Table 1 and Fig. 3, indicating the increase of X-ray density of  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$  with increasing  $\text{Bi}^{3+}$ .

Table 1

The structural and magnetic characteristics (lattice parameter,  $a$ , X-ray density,  $d_x$ , saturation magnetization,  $M_s$ , Néel temperature,  $T_N$ , and the anisotropy constant,  $K_1$ ) of  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$  ( $x=0, 0.1, 0.15$ ) samples. The  $M_s$  values of the as-prepared bulk ferrites are compared with those of the corresponding nanocrystalline  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$  counterparts [9].

$x$	$a$ (Å)	$d_x$ ( $\text{g cm}^{-3}$ )	$M_s$ ( $\text{emu g}^{-1}$ )	$M_s$ ( $\text{emu g}^{-1}$ ) <sup>a</sup>	$T_N$ (K)	$K_1$ ( $\text{erg cm}^{-3}$ )
0	8.334	5.379	44	42	848	76723
0.1	8.341	5.716	40	35	788	51547
0.15	8.346	5.881	38	29	753	50284

<sup>a</sup>Data taken from Ref. [9].

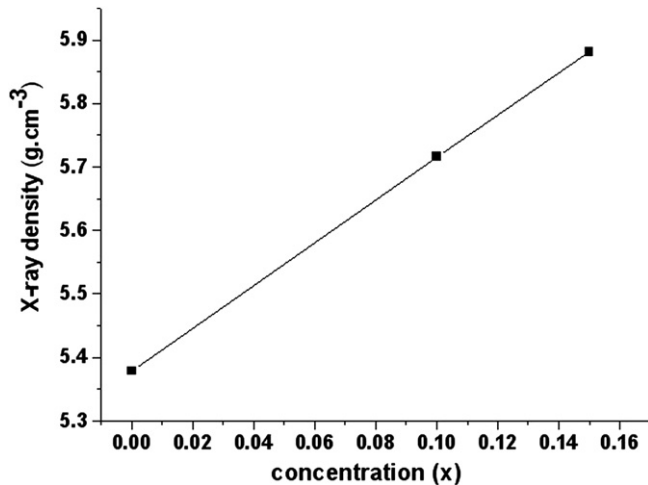


Fig. 3. The X-ray density of  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$  versus Bi content ( $x$ ).

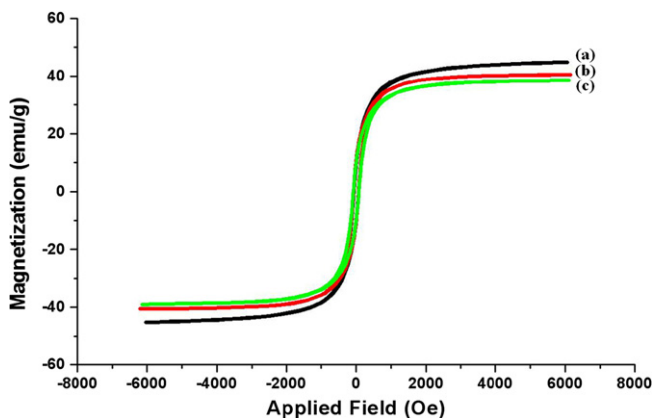


Fig. 4. The hysteresis loops for  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$  samples: (a)  $x=0$ , (b)  $x=0.1$ , (c)  $x=0.15$ , recorded at room temperature.

Fig. 4 shows the hysteresis curves of the  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$  ( $x=0, 0.1, 0.15$ ) samples. The measured values of the saturation magnetization ( $M_s$ ) for all the samples are presented in Table 1, indicating the decrease of  $M_s$  with increasing  $\text{Bi}^{3+}$  concentration. This behavior is explained by the magnetic dilution of the ferrite system with nonmagnetic  $\text{Bi}^{3+}$  ions; i.e., in the present case, an  $\text{Fe}^{3+}$  ion with an effective magnetic moment of  $5\mu_B$  is substituted by a nonmagnetic  $\text{Bi}^{3+}$  cation. Note that a very

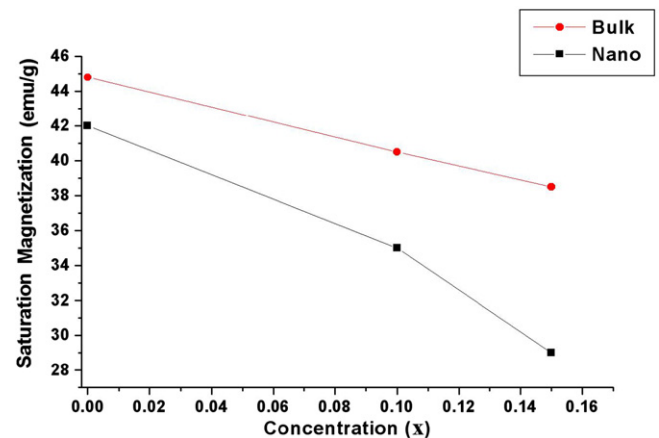


Fig. 5. Magnetization of bulk and nanocrystalline  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$  samples versus  $x$ . The  $M_s$  values of nanocrystalline  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$  are taken from Ref. [9].

small hysteresis for all the as-prepared samples indicates that the Ni–Bi ferrite can be used most efficiently as a soft magnetic material with a negligible hysteresis loss [21]. From the qualitative point of view, these results are similar to those reported in our work for nanocrystalline Ni–Bi ferrite prepared via the sol–gel method [9]. Quantitatively, the saturation magnetization for the bulk  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$  samples is significantly greater than that of corresponding nanocrystalline counterparts, see Fig. 5. As it is generally accepted, a reduction of the saturation magnetization with decreasing size of ferrite particles can be attributed to the presence of canted spins and nonequilibrium cation distribution in the surface shell of nanoparticles [3,6,22,23].

The temperature-dependent magnetic measurements (Fig. 6) revealed that the Néel temperature ( $T_N$ ) of the as-prepared  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$  shifts towards lower temperature with increasing  $\text{Bi}^{3+}$  concentration;  $T_N=848, 788$  and  $753$  K for  $x=0, 0.1$  and  $0.15$ , respectively. The Néel temperature of a ferrimagnetic spinel oxide has been found to be closely related to the number of  $\text{Fe}^{3+}(\text{A})\text{--O}^{2-}\text{--Fe}^{3+}[\text{B}]$  linkages per formula unit of a material and also to the distribution of  $\text{Fe}^{3+}$  ions over the (A) and [B] sites [24]. Thus, the observed reduction of  $T_N$  with increasing  $x$  in  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$  may be interpreted as a consequence of weakening of the (A)–[B] superexchange interaction due to

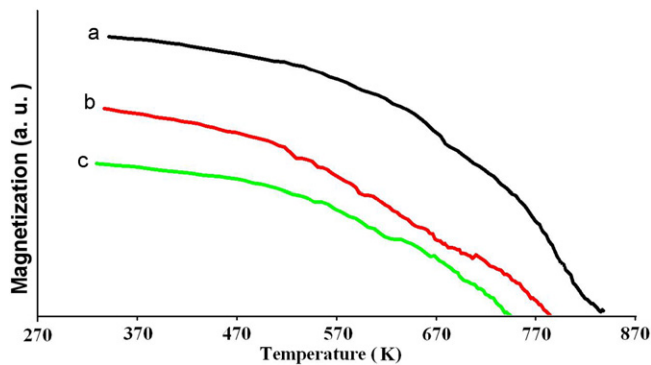


Fig. 6. The temperature-dependent magnetization for  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$  samples: (a)  $x=0$ , (b)  $x=0.1$ , (c)  $x=0.15$ , recorded at  $H=0.002$  T.

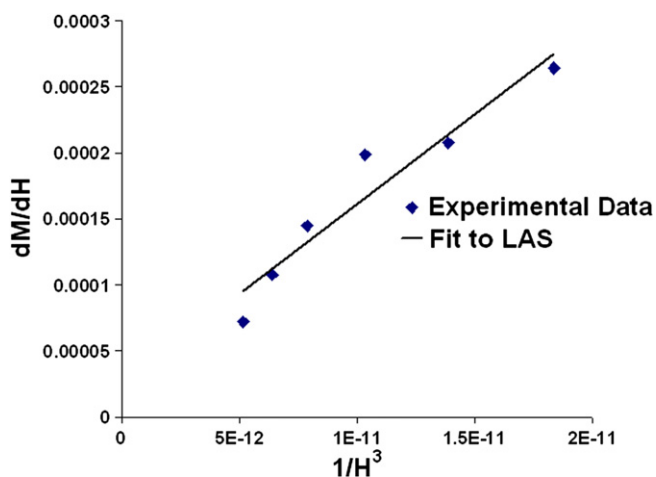


Fig. 7. Magnetization data for  $\text{NiFe}_{1.9}\text{Bi}_{0.1}\text{O}_4$  fitted using the LAS model.

the substitution of the  $\text{Fe}^{3+}$  ions by nonmagnetic  $\text{Bi}^{3+}$  cations. In such a case, the internal energy required to offset the spin alignment will decrease and the Néel temperature is expected to fall.

The anisotropy constant for the  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$  samples ( $x=0, 0.1, 0.15$ ) was estimated using the “Law of approach to saturation (LAS)” model [24]. According to this theory, the dependence of magnetization,  $M$ , on the applied magnetic field,  $H$ , is given by

$$M = M_s[1 - (b/H^2)] + kH \quad (3)$$

where  $b=(8/105) (K_1^2/\mu_0^2 M_s^2)$ ,  $\mu_0$  is the permeability of the free space,  $K_1$  is the cubic anisotropy constant and the term  $kH$  is known as the forced magnetization. Note that the forced magnetization causes a linear increase in the spontaneous magnetization especially at high fields. The numerical coefficient 8/105 applies to cubic anisotropy of random polycrystalline samples. In the present series, the experimental data [ $M$ – $H$  curve] above 0.5 T (high field parts) are fitted to Eq. (3) (with the fitting parameters  $M_s$  and  $K_1$ ) neglecting the forced magnetization term. A typical curve (straight line) fitted to LAS is shown in Fig. 7 for the sample  $\text{NiFe}_{1.9}\text{Bi}_{0.1}\text{O}_4$ . The anisotropy

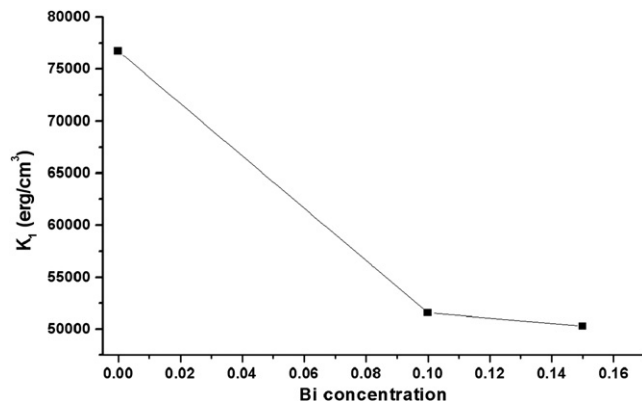


Fig. 8. Variation of the cubic anisotropy constant with  $x$  for  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$ .

constants of the  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$  samples ( $x=0, 0.1$  and  $0.15$ ) with various  $\text{Bi}^{3+}$  concentrations are shown in Fig. 8. As it is seen, an increasing  $\text{Bi}^{3+}$  concentration leads to a decrease in the anisotropy constant. In pure  $\text{NiFe}_2\text{O}_4$ , the  $\text{Fe}^{3+}$  ions occupy both (A) and [B] sites; thus the strong anisotropy is primarily due to the presence of  $\text{Fe}^{3+}$  ions on the octahedral sites of the spinel structure. The crystal field (trigonal field) is not capable of removing the orbital degeneracy of  $\text{Fe}^{3+}$  ions at the octahedral sites so that the orbital magnetic moment is not quenched and therefore there is a strong spin–orbit coupling, which produces large magnetocrystalline anisotropy energy. The substitution of  $\text{Fe}^{3+}$  ions by  $\text{Bi}^{3+}$  cations on [B] sites in  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$  [9] reduces the spin–orbit coupling, which may lead to a decrease in the anisotropy constant of the substituted ferrites.

#### 4. Conclusions

Nickel–bismuth ferrites in the bulk form with various  $\text{Bi}^{3+}$  concentrations ( $x=0, 0.1, 0.15$ ) were synthesized by the thermal treatment of  $\text{NiFe}_{2-x}\text{Bi}_x\text{O}_4$  nanoparticles previously prepared by sol–gel technique. The lattice parameter of the as-prepared cubic ferrites increases with increasing  $x$  due to the larger ionic radius of  $\text{Bi}^{3+}$  in comparison to that of  $\text{Fe}^{3+}$  ions. The decrease in the saturation magnetization and Néel temperature of the samples with increasing  $\text{Bi}^{3+}$  concentration is explained by the effect of magnetic dilution of the ferrite system and, consequently, by the weakening of the (A)–[B] super-exchange interaction, respectively. The substitution of the magnetic  $\text{Fe}^{3+}$  ions by nonmagnetic  $\text{Bi}^{3+}$  cations reduces the spin–orbit coupling leading to a decrease in the anisotropy constant of the ferrites. The Bi-substituted nickel ferrites exhibit a soft magnetic behavior.

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