

# Improving ferroelectric and piezoelectric properties of $\text{PbFe}_{1/4}\text{Sc}_{1/4}\text{Nb}_{1/2}\text{O}_3$ ceramics by oxide doping prepared via a B-site oxide mixing route

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## Abstract

Oxide-doped  $\text{PbFe}_{1/4}\text{Sc}_{1/4}\text{Nb}_{1/2}\text{O}_3$  (PFSN) ceramics were prepared by the conventional ceramic processing via a B-site oxide mixing route. The oxide-doped PFSN ceramics exhibit pure rhombohedral perovskite structure and high relative density. Due to the  $\text{MnO}_2$  and  $\text{Li}_2\text{CO}_3$  doping, the synthesized ceramics exhibit improved dielectric property, and excellent ferroelectric and piezoelectric properties, which can be attributed to a charge compensation mechanism where donor defects are generated by different ionization paths. The values of remnant polarization  $P_r$ , electromechanical coupling coefficient  $K_p$  and piezoelectric constant  $d_{33}$  of the  $\text{MnO}_2$ - and  $\text{Li}_2\text{CO}_3$ -doped PFSN ceramics are  $16.35 \mu\text{C}/\text{cm}^2$ , 0.232 and  $112.7 \text{ pC}/\text{N}$ , and  $29.22 \mu\text{C}/\text{cm}^2$ , 0.241 and  $112.5 \text{ pC}/\text{N}$ , respectively. Therefore, oxide doping is an efficient method to improve dielectric, ferroelectric and piezoelectric properties of the PFSN-based ferroelectric ceramics. © 2012 Elsevier Ltd and Techna Group S.r.l. All rights reserved.

**Keywords:** C. Ferroelectric property; C. Piezoelectric property; Crystal structure;  $\text{PbFe}_{1/4}\text{Sc}_{1/4}\text{Nb}_{1/2}\text{O}_3$ -based ceramics

## 1. Introduction

Relaxor ferroelectric ceramics exhibit excellent electrical properties, which make them promising applications in multilayer capacitor and related versatile application fields [1–5]. 1:1-type lead-based complex perovskite relaxors  $\text{Pb}(\text{Fe}_{1/2}\text{Nb}_{1/2})\text{O}_3$  (PFN) and  $\text{Pb}(\text{Sc}_{1/2}\text{Nb}_{1/2})\text{O}_3$  (PSN) attract our research attention since PFN and PSN exhibit rhombohedral perovskite structure and complete solid solution  $(1-x)\text{Pb}(\text{Fe}_{1/2}\text{Nb}_{1/2})\text{O}_3-x\text{Pb}(\text{Sc}_{1/2}\text{Nb}_{1/2})\text{O}_3$  (PFN–PSN) can form in the whole composition range [6–9].

The nature of relaxor behavior has attracted enormous fundamental investigations. Usually, the relaxor behavior is attributed to the polar clusters and/or the nanoscopic inhomogeneity induced by the random occupation of the crystallographic equivalent positions by different cations [10,11]. Such structural character produces short-range occupational ordering and engineered domain configuration,

which prevents the complete development of long-range ferroelectric ordering, destroys the normal ferroelectric phase transition and makes the physical properties of the relaxor ferroelectrics similar to those of the disordered ferromagnetic materials [12,13]. However, the understanding of the relaxor properties is not yet satisfactory [14,15].

In this paper, the composition  $\text{PbFe}_{1/4}\text{Sc}_{1/4}\text{Nb}_{1/2}\text{O}_3$  (PFSN) is adopted for further optimization since PFSN crystallizes into a pseudo-cubic structure around room temperature and exhibits relatively large dielectric constant, low sintering temperature, and relaxor behavior. Since the iron- and scandium-containing ferroelectric ceramics tend to exhibit large dielectric loss and strong dielectric frequency dispersion, the pure PFSN ceramics exhibit deteriorated ferroelectric and piezoelectric properties. In views of crystal chemistry and materials engineering, the differences of ionic size, atomic weight, electric negative, and/or charge are produced by cation substitution on the crystallographic equivalent A-site and/or B-site of the perovskite structure [16,17], which provides straightforward experimental method for systematically tuning electrical properties of piezoelectric

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materials [18]. Therefore, chemical doping  $\text{MnO}_2$  and  $\text{Li}_2\text{CO}_3$  are used to improve electrical properties of the Fe- and Sc-containing ferroelectric ceramics. Through adjustment of the composition of the solid solution, defect structure, spontaneous polarization, electrostatic interactions, and ordering degree of the perovskite structure may be influenced, which provides the possibility to tailor electrical properties of PFN–PSN [7,16,17,19].

## 2. Experimental procedure

The oxide-doped PFSN ceramics were prepared by the conventional solid-state reaction method via a B-site oxide mixing route, in which all oxides of the B-site of the perovskite structure were pre-calcined simultaneously instead of preparing different wolframite precursors separately [20]. High-purity oxides and carbonate,  $\text{PbO}$  (> 99.9%),  $\text{Fe}_2\text{O}_3$  (> 99.99%),  $\text{Sc}_2\text{O}_3$  (> 99.9%),  $\text{Nb}_2\text{O}_5$  (> 99.9%),  $\text{MnO}_2$  (> 99.95%), and  $\text{Li}_2\text{CO}_3$  (> 99.9%), were used as raw materials. To maintain stoichiometry the raw materials were dried before weighing and the synthesized B-site precursors were weighed and introduced into the batch weighing. The well-mixed stoichiometric B-site oxides,  $\text{Fe}_2\text{O}_3$ ,  $\text{Sc}_2\text{O}_3$  and  $\text{Nb}_2\text{O}_5$ , were pre-calcined at 1200 °C for 3 h. After crushing and grinding, stoichiometric  $\text{PbO}$  and 0.25 wt% oxide dopants were added to the synthesized B-site precursors, and the mixed powders were calcined at 900 °C for 2 h. X-ray diffraction measurement (XRD, Rigaku RINT-2200 VS diffractometer) confirmed that the calcined powders exhibit pure perovskite structure. Then the calcined powders were isostatically pressed into pellets at a pressure of 400 kgf/cm<sup>2</sup> with the addition of 1 wt% polyvinyl alcohol (PVA) and sintered at 1160 °C for 2 h. The green pellets were buried under a mixture of raw materials with equal weight and same composition of the pellets in a covered alumina crucible to minimize the evaporation of lead during sintering.

The sintered oxide-doped PFSN ceramics were ground and polished to obtain flat and parallel surfaces. XRD measurement confirmed that the synthesized ceramics maintain perovskite structure during the sintering process. The elemental distribution of the sintered ceramics was measured by an Energy dispersive X-ray fluorescence elemental analyzer (XRFEA, RoHS SEA-2010). For electrical properties characterization, silver paste was coated on both surfaces of the well-polished ceramics and fired at 650 °C for 15 min to provide robust electrodes. Dielectric property under direct current (DC) bias voltage was measured by a Turnkey broadband dielectric spectrometers Concept 42 (Novocontrol Technologies GmbH & Co. KG, Germany). Polarization–electric field ( $P$ – $E$ ) hysteresis loop was characterized by a Radiant Precision Premier LC ferroelectric material test system (Radiant Technologies Inc., USA), in which a Sigma Model M10 chamber was attached to measure the  $P$ – $E$  loop upon heating. The Sigma Model M10 chamber was controlled by a Sigma Systems Model C4 programmable temperature

controller, which can measure the thermal range –100 to 200 °C (Sigma Systems Corporation, USA). During the  $P$ – $E$  hysteresis loop measurement, the ceramics were immersed in silicon oil to prevent arching. For piezoelectric property measurement, the oxide-doped PFSN ceramics were poled under an electric field of 2 kV/mm at 80 °C for 15 min in silicon oil and then slowly cooled down to room temperature with maintaining half of the applied electric field. Piezoelectric property was measured by a ZJ-6A Berlincourt-type quasi-static  $d_{33}/d_{31}$  meter (Institute of Acoustics, Chinese Academy of Sciences, Beijing, China). Detailed performance measurement procedures were described elsewhere [21].

## 3. Results and discussion

In previous work [22], we reported that the 1160 °C sintered oxide-doped PFSN ceramics exhibit pure rhombohedral perovskite structure, densified microstructure morphology, high relative density and improved dielectric property, which are especially favorable for improving ferroelectric and piezoelectric properties of the PFSN-based ceramics.

The elemental distribution is important for understanding of the effects of oxide doping on the electrical properties of the PFSN-based ceramics. Fig. 1 shows XRFEA spectra of the  $\text{MnO}_2$ -doped PFSN ceramics. The average values of elemental contents based on the XRFEA measurement are shown in the figure. The elemental contents at different sites exhibit slight difference, especially for  $\text{Mn}^{2+}$ , indicating that the preparation of homogeneous ceramics is difficult, which will affect local crystal structure and electrical properties of the sintered ceramics. However, the variation of elemental concentration is not apparent. Therefore, the nominal composition can be used to represent the real composition of the oxide-doped PFSN ceramics within the detection sensitivity of XRFEA.

The shape of  $P$ – $E$  ferroelectric hysteresis loops is influenced greatly by electric field strength and frequency. Fig. 2

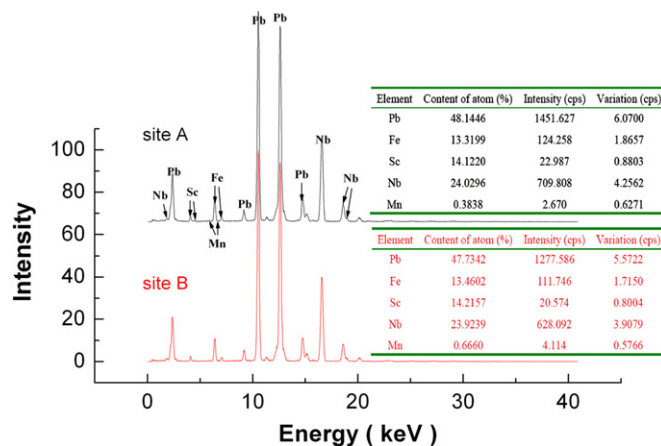


Fig. 1. Energy dispersive X-ray fluorescence spectra of the  $\text{MnO}_2$ -doped PFSN ceramics measured at different sites. Inset of the figure shows the average values of elemental content.

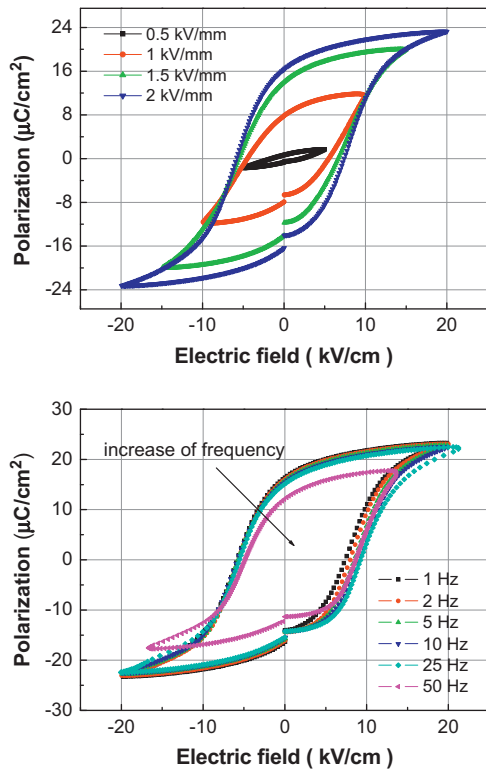


Fig. 2. Room-temperature  $P$ – $E$  hysteresis loops of the  $\text{MnO}_2$ -doped PFSN ceramics measured at different electric fields and frequencies.

shows  $P$ – $E$  hysteresis loops of the  $\text{MnO}_2$ -doped PFSN ceramics. For ferroelectric ceramics, large electric field can afford enough energy for ferroelectric domains to switch and align along the direction of the external electric field. Therefore, at 0.5 kV/mm, a nearly linear  $P$ – $E$  relationship is observed, and at 2 kV/mm, a saturate  $P$ – $E$  loop is fully developed, correlating with the activation energy needed for the rotation of the ferroelectric domains. For all the electric field measured, the obtained  $P$ – $E$  hysteresis loops are rather symmetric and no apparent pinning effect appears. Low frequency is favorable for different polarization mechanisms to follow up the change of the external electric field. Between 1–25 Hz, the  $P$ – $E$  hysteresis loops are fully developed, saturated and symmetric with almost the same values of remnant polarization  $P_r$  and coercive field  $E_C$ . At 50 Hz, the  $P$ – $E$  hysteresis loop exhibits apparent asymmetry, indicating that the pinning-down phenomenon exists and the measurement frequency is too fast for the ferroelectric domains to rotate with the electric field.

Detailed room-temperature  $P$ – $E$  hysteresis loops of the oxide-doped PFSN ceramics are shown in Fig. 3. Both the  $\text{MnO}_2$ - and  $\text{Li}_2\text{CO}_3$ -doped PFSN ceramics exhibit fully developed, symmetric and saturated  $P$ – $E$  hysteresis loops. The hysteresis loops are sharp and narrow accompanied by small value of  $E_C$ , which is similar to that of rhombohedral perovskite structure ferroelectrics. Based on the fully saturated  $P$ – $E$  hysteresis loops, ferroelectric properties  $P_r$  and  $E_C$  can be determined, which are shown in Fig. 3. As compared to the  $\text{MnO}_2$ -doped PFSN ceramics, the  $\text{Li}_2\text{CO}_3$ -doped

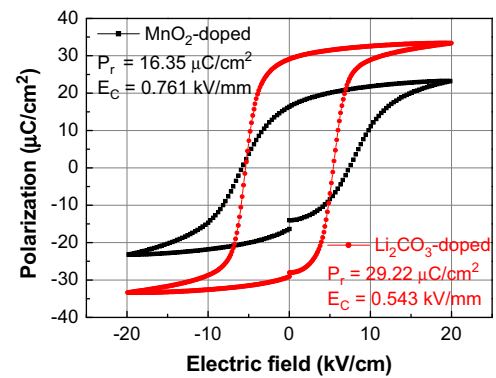


Fig. 3. Room-temperature  $P$ – $E$  hysteresis loops of the oxide-doped PFSN ceramics measured at 1 Hz and 2 kV/mm.

PFSN ceramics exhibit large value of  $P_r$  and small value of  $E_C$ , being  $29.22 \mu\text{C}/\text{cm}^2$  and  $0.543 \text{ kV}/\text{mm}$ , respectively, which is considered as correlating with different charge compensation mechanism induced by oxide doping [22].

Fig. 4 shows room-temperature 3-D graph of relative dielectric constant-frequency-DC bias voltage of the oxide-doped PFSN ceramics. The dielectric response peaks appearing at different frequencies can be attributed to the different vibration modes of electromechanical coupling of the poled ceramics. The dielectric response character is influenced greatly by oxide doping, in which the value of dielectric constant and the shape of dielectric resonant response spectra exhibit great differences. DC bias electric voltage also influences the dielectric response character apparently, which is considered as relating to the electric field induced micro-macro ferroelectric domain change and/or the structural ferroelectric phase transition. Usually, such mechanism can be confirmed by the dielectric-temperature spectrum measurement using the poled ceramics. After poling, an additional dielectric peak or shoulder will be induced, which is considered as correlating with the orientation of ferroelectric domain or the macro-micro domain change upon heating [21]. However, for the oxide-doped PFSN ceramics, the additional electric field-induced minor dielectric peak or shoulder may move to cryogenic temperature, which cannot be measured by our experimental.

For piezoelectric applications the electromechanical coupling coefficient is an important parameter for piezoelectric materials. Frequency dependence of impedance and phase degree of the  $\text{MnO}_2$ -doped PFSN ceramics is shown in Fig. 5. The  $\text{MnO}_2$ -doped PFSN ceramics exhibit perfect resonant response peak, which provides efficient energy conversion between electrical energy and mechanical energy. The electromechanical coupling coefficient  $K_p$  and the mechanical equality factor  $Q_m$  of the  $\text{MnO}_2$ -doped PFSN ceramics can be calculated from the equations  $K_p \approx \sqrt{\frac{f_p - f_s}{f_s}} \times 2.51$  and  $Q_m = \frac{f_p^2}{2\pi f_s C^T R(f_p^2 - f_s^2)}$ , being 0.232 and 203.8, respectively, at room-temperature.

Environmental temperature exerts great influence on piezoelectric properties of the PFSN-based ceramics due to the existence of depolarization phenomenon of piezoelectric

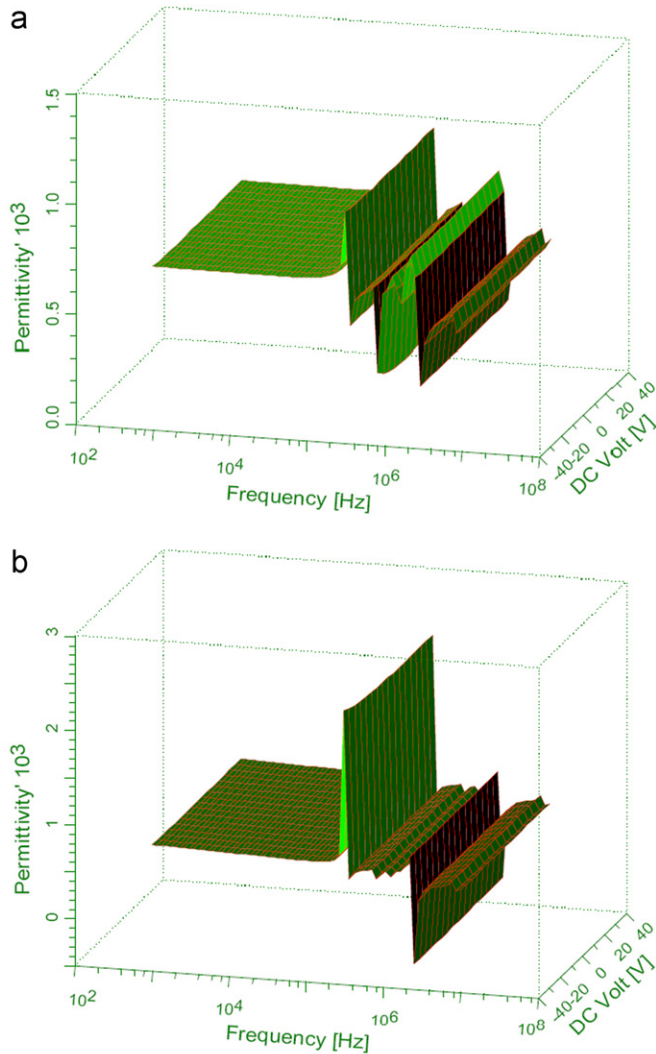


Fig. 4. Frequency dependence of dielectric constant under different DC bias voltage of the poled (a)  $\text{MnO}_2$ -doped PFSN ceramics and (b)  $\text{Li}_2\text{CO}_3$ -doped PFSN ceramics measured at room-temperature.

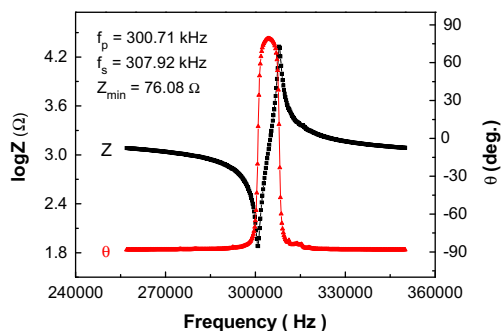


Fig. 5. Frequency dependence of impedance and phase angle degree in radial-extension-vibration-mode of the  $\text{MnO}_2$ -doped PFSN ceramics measured at room-temperature.

materials. Temperature dependence of radial-extension-vibration electromechanical coupling coefficient  $K_p$  and piezoelectric constant  $d_{33}$  of the oxide-doped PFSN ceramics

is shown in Figs. 6 and 7. These curves can nearly be divided into two regions, in which the values of  $K_p$  and  $d_{33}$  decrease with different slope upon temperature. The values of  $K_p$  and  $d_{33}$  decrease slightly from room-temperature to 80 °C, while decrease greatly after 80 °C, till the end of temperature measured. The dramatic decrease of the values of  $K_p$  and  $d_{33}$  is considered as correlating with the ferroelectric phase transition from rhombohedral ferroelectric phase to cubic paraelectric phase, and can be measured by dielectric property-temperature spectrum. However, the temperature that the values of  $K_p$  and  $d_{33}$  decreases greatly is rather below than the temperature of the dielectric constant maximum  $T_m$  determined by dielectric measurement, which is usually called the depolarization temperature  $T_d$  [23]. Therefore, the service temperature is determined based on the experimental results, being around 80 °C, and higher temperature of  $T_d$  can be obtained by the addition of  $\text{PbTiO}_3$  in the PFSN solid solution [9].

Based on the above experimental results and the results reported before, we can draw a conclusion. Oxide doping is an efficient method to improve electrical properties of the PFSN-based ferroelectric ceramics, in which not only the dielectric response peaks become sharp, the dielectric loss decreases, and the dielectric frequency dispersion is suppressed greatly, but also the ferroelectric and piezoelectric properties are improved. Such phenomena can be interpreted

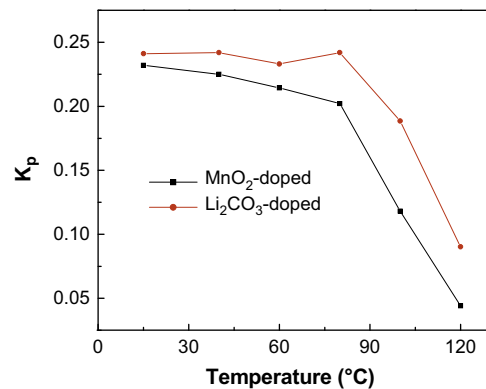


Fig. 6. Temperature dependence of radial electromechanical coupling coefficient  $K_p$  of the oxide-doped PFSN ceramics.

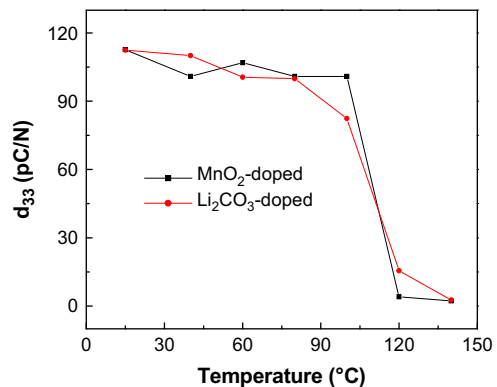


Fig. 7. Temperature dependence of piezoelectric constant  $d_{33}$  of the oxide-doped PFSN ceramics.



by a charge compensation mechanism, in which donor defects are generated by manganese ionization  $Fe_{Fe}^x/Sc_{Sc}^x + Mn(3d^5 4s^2) + E_D \rightarrow Mn_{Fe/Sc} + e'$  and lithium sitting at interstitial sites of the perovskite structure  $V_i + Li(2s^1) + E_D \rightarrow Li_i + e'$ , respectively.  $Li_2CO_3$  doping exhibits enhanced improvement of electrical properties of the PFSN-based ceramics, which can be attributed to the small ionization energy of the interstitial ionic defects, leading to superior compensating for charge disproportion. Therefore,  $MnO_2$  and  $Li_2CO_3$  doping exhibit significant meaning in improving ferroelectric and piezoelectric properties of the PFN- and PSN-based ceramics.

#### 4. Conclusions

Pure rhombohedral perovskite structure oxide-doped PFSN ceramics were prepared by the conventional solid-state reaction method via the B-site oxide mixing route. The oxide-doped PFSN ceramics exhibit large relative density and rather homogenous microstructure morphology. Due to the  $MnO_2$  and  $Li_2CO_3$  doping, the synthesized ceramics exhibit improved dielectric property, and excellent ferroelectric and piezoelectric properties, which are considered as relating to the charge compensation mechanism where donor defects are produced by different ionization paths. The values of  $P_r$  and  $E_C$  of the  $MnO_2$ - and  $Li_2CO_3$ -doped PFSN ceramics are  $16.35 \mu C/cm^2$  and  $0.761 kV/mm$ , and  $29.22 \mu C/cm^2$  and  $0.543 kV/mm$ , respectively. The values of  $K_p$ ,  $Q_m$  and  $d_{33}$  of the  $MnO_2$ - and  $Li_2CO_3$ -doped PFSN ceramics are 0.232, 203.8 and  $112.7 pC/N$ , and 0.241, 320.7 and  $112.5 pC/N$ , respectively. The values of  $K_p$  and  $d_{33}$  decrease greatly after  $80^\circ C$ , which is considered as correlating with the ferroelectric phase transition from rhombohedral ferroelectric phase to cubic paraelectric phase. Therefore, oxide doping is an efficient method to improve dielectric, ferroelectric and piezoelectric properties of the PFSN-based ferroelectric ceramics.

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