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X-ray diffraction, Raman and photoluminescence studies of nanocrystalline cerium oxide thin films

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Abstract

Cerium oxide (CeO₂) thin films were deposited on Si (100) and glass substrates at various oxygen partial pressures (2×10^{-5} mbar- 3×10^{-1} mbar) and a substrate temperature of 673 K, by pulsed laser deposition (PLD). The structural, morphological and optical properties of the films were characterized by X-ray diffraction, Raman spectroscopy, high resolution transmission electron microscopy, atomic force microscopy, photoluminescence and UV-visible spectroscopy. XRD analysis revealed the polycrystalline and cubic structure of the CeO₂ films with preferred orientation for (200) as the increase of oxygen partial pressure. The Raman studies indicated the formation of Ce–O with the systematic variation of peak intensity and FWHM with oxygen partial pressures. The high resolution transmission electron microscopy investigation confirmed the polycrystalline and cubic nature of the CeO₂ films with (200) preferred orientation. The AFM studies showed the roughness (RMS) values of the films increased from 0.8 nm to 4.6 nm with increasing oxygen partial pressure from 2×10^{-5} mbar to 3×10^{-1} mbar. The photoluminescence (PL) investigation indicated the bandgap values in the range 3.05–3.10 eV with increasing oxygen partial pressures. The UV-visible spectroscopy analysis demonstrated that the refractive index of the films decreased from 2.41 to 1.72 with increasing oxygen partial pressures.

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1. Introduction

Cerium oxide (CeO₂) thin film is of great interest due to its excellent properties such as high refractive index, wide bandgap, high dielectric constant, high melting point, high transparency in the VIS–NIR regions with high thermal and chemical stability [1,2]. These properties enhance the use of CeO₂ films for various electronic applications, such as siliconon-insulator structure, miniaturized stable capacitors and buffer layers for the high temperature superconductor films [3,4].

CeO₂ films are also used in single and multilayer coatings for optical devices, electro-chromic windows, sensors and high temperature oxidation resistant coatings [5–7]. It is widely used as UV blocking filters in medical glassware and aerospace windows due to strong absorption in ultraviolet region [8]. High quality CeO₂ films have interesting optical properties for application as counter electrodes in transparent opto-ionic devices. It is a promising material for fast oxygen sensors at high temperature, because of its chemical stability and high diffusion coefficient of oxygen vacancies. Doped and undoped CeO₂ are important luminescent materials and investigated for optical marker, biological labeling and phosphor applications [9,10]. CeO₂ thin films are prepared by several methods, such as spray pyrolysis [11], electron-beam evaporation [12], sputtering [13], electro deposition [14] and pulsed

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laser deposition (PLD) [15,16]. Among these methods, PLD is a simple and unique technique to prepare high quality thin films of metals, semiconductors and insulators [16]. One of the main advantages of PLD is the generation of hyperthermal species with high kinetic energy of the order of 100 eV [17]. Deposition of hyperthermal species can enhance the adatom mobility and hence the film quality. In the PLD process, there are several operating parameters such as wavelength, fluence, repetition rate, energy per pulse, ambient gases, pressure of the ambient gases, target to substrate distance and substrate temperature. Among these parameters, oxygen partial pressure and substrate temperature have profound influence on the properties. Hence, the present work deals with the influence of oxygen partial pressure on microstructural and optical properties of CeO₂ thin films.

2. Experimental details

CeO₂ (99.99% purity) powder was compacted into a pellet of 25 mm diameter and 4 mm thickness using a uni-axial press. The pellet was sintered at 1673 K for 6 h and used as target for preparation of thin films. Deposition of CeO₂ films on clean Si (100) and glass substrates was carried out as a function of oxygen partial pressure $(2.0 \times 10^{-5} - 3.0 \times 10^{-1} \text{ mbar})$ at 673 K, using KrF (248 nm) excimer laser with the energy density of 3 J/cm². The other deposition parameters are given elsewhere [15]. The structural property of the films was studied by X-ray diffractometer (X'pert PW 3040 D-8, PANalytical) using Cu K₀₁ radiation.

The crystalline sizes were determined from the XRD data using Scherrer formula:

$$D = \frac{K\lambda}{\beta Cos\theta}$$

where $\beta = \sqrt{(B^2 - b^2)}$ K is Scherrer constant (0.9), λ is the wavelength of the $CuK_{\alpha 1}$ radiation, D is the crystallite size, β is full width at half maximum (FWHM) after correcting for instrumental broadening, θ is the diffraction angle, B is the observed FWHM of the film and b is the instrumental broadening. The structural property was also examined by Raman spectrophotometer (NRS-3300) at room temperature and ambient air. The Nd-YAG laser of 7.7 mW power with a wavelength of 532 nm was used as the excitation source. A high resolution transmission electron microscope (HRTEM) study of the CeO₂ film was carried out by a JEOL 2000 EX II (T) transmission electron microscope operated at 200 kV. The surface morphology and RMS surface roughness of the films was analyzed by atomic force microscope (AFM) (XE-100 Park systems). Room temperature photoluminescence (PL) was recorded using T64000 Jovin-Yvon spectrometer with a CCD detector and 1800 g/mm grating. A He-Cd laser of wavelength 325 nm, was used as the excitation source. The optical properties was studied by UV-VIS-NIR (Model no.: 3101/PC, Shimadzu) spectrophotometer in the wavelength range 190-800 nm. The refractive indices of films were calculated from transmittance spectra using the following equation:

$$n = [N + (N^2 - n_0^2 n_1^2)^{1/2}]^{1/2}$$

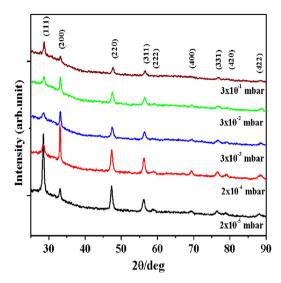


Fig. 1. XRD pattern of ceria films deposited on glass substrates at various oxygen partial pressures and 673 K.

Table 1
Oxygen partial pressure versus crystallite size, Raman FWHM and refractive index.

Sl. no.	Oxygen partial pressure (mbar)	Crystallite size (nm) (111)	Crystallite size (nm) (200)	FWHM from Raman (cm ⁻¹)	Refractive index at 550 nm
1	2.0×10^{-5}	25	24	12.54	2.41
2	2.0×10^{-4}	12	45	15.29	2.34
3	3.0×10^{-3}	7	28	10.10	1.91
4	3.0×10^{-2}	11	24	9.79	1.77
5	3.0×10^{-1}	32	19	_	1.72

where

$$N = \frac{n_0^2 + n_1^2}{2} + 2n_0n_1 \frac{T_{max} - T_{min}}{T_{max}T_{min}}$$

where n is the refractive index of the film, n_o is the refractive index of air, n_I is the refractive index of the substrate, T_{max} , T_{min} are the maximum and minimum transmittance respectively.

3. Results and discussion

3.1. Microstructural characterization

3.1.1. X-ray diffraction studies

Fig. 1 shows the XRD pattern of the CeO₂ films deposited on glass substrates at various oxygen partial pressures. The films showed (111), (200), (220), (311), (400), (331), (420) and (422) reflections correspond to the polycrystalline cubic structure of CeO₂ [18,19]. The films showed higher intensity for (111) reflection at base vacuum of 2.0×10^{-5} mbar. There was an increased intensity for (200) reflection at higher oxygen partial pressure $(2 \times 10^{-4} - 3.0 \times 10^{-2} \text{ mbar})$, whereas the intensity of (111) reflection decreased simultaneously, which indicates the preferred orientation of the films with increasing oxygen partial pressure [20,21]. The thickness of the films was measured and

found to decrease from 300 nm to 70 nm with increasing oxygen partial pressures. This is due to the increased scattering of ablated species with oxygen molecules in the deposition chamber. The peak intensity of all the reflections decreased at an oxygen partial pressure of 3×10^{-1} mbar, due to the reduced thickness. The crystallite size was calculated from Scherrer's equation after subtracting the instrumental broadening (Table 1). Kanakaraju et al. [21] investigated the properties of CeO₂ films prepared on fused silica substrates as a function of oxygen partial pressure and observed the texture of (200) in the temperature ≥400 °C. Amirhaghi et al. [22], grown the CeO₂ films on Si (100) substrates at different oxygen partial pressure and temperature using PLD and revealed that the films deposited at lower oxygen partial pressure, high substrate temperature and the laser fluence $> 2 \text{ J cm}^{-2}$ were good quality films. Wang et al. [20] studied the CeO₂ films deposited on Si (100) and found that the films orientation varied with increasing oxygen pressure.

3.1.2. Raman spectroscopy studies

The Raman spectroscopy was used to study the structure of the films deposited on Si (100). In general, CeO2 has a fluorite structure and six optical phonon modes, which gives three zone center frequencies, 272, 465 and 595 cm⁻¹ correspond to doubly degenerate transverse optical (TO) mode, triply degenerate and the non-degenerate longitudinal optical (LO) mode. CeO₂ possess cubic structure with a symmetrical stretching mode of the Ce-O vibrational unit. Fig. 2a shows the Raman peak of sintered CeO₂ pellet, appeared at 463 cm⁻¹ and is very sharp and symmetric. Fig. 2b shows the Raman spectra of CeO₂ thin films deposited in the oxygen partial pressure range 2.0×10^{-5} – 3.0×10^{-1} mbar. All the films showed peak position $\sim 463 \text{ cm}^{-1}$ and is associated with the F_{2g} Raman active mode of polycrystalline cubic structure of CeO₂. The peak $\sim 300 \text{ cm}^{-1}$ is due to Si (100) substrate. At base pressure $(2.0 \times 10^{-5} \text{ mbar})$, the Raman active mode appeared \sim 462 cm⁻¹ with higher peak intensity and lower FWHM of 13 cm⁻¹. The frequency shift was observed in the range 462- 466 cm^{-1} , as the oxygen partial pressure increased from 2×10^{-4}

to 3.0×10^{-2} mbar. The FWHM decreased with increasing oxygen partial pressure [23–25]. No peak was observed in the Raman spectrum at 3.0×10^{-1} mbar, due to poor crystallinity and these results are in accordance with XRD results (Fig. 1). The spectra are also prone to be directly affected by the presence of oxygen vacancies. The peak shape becomes broader and asymmetric progressively with increasing oxygen partial pressure. Raman peak broadening indicates the formation of nanocrystalline phase of CeO₂. The Raman study shows the systematic variation in the peak intensity, FWHM and asymmetry with the increase of oxygen partial pressure and is consistent with XRD results. The similar results were also observed by Jonathan et al. [25], Kanakaraju et al. [21] and Wang et al. [23].

3.1.3. High resolution transmission electron microscopy (HRTEM) analysis

The microstructure of CeO₂ film deposited on NaCl single crystal at 3×10^{-2} mbar of oxygen partial pressure and 673 K was analyzed by high resolution transmission electron microscope (HRTEM). Fig. 3a shows the low magnification bright field image of CeO₂ and uniform crystallites distribution. Fig. 3b shows the selected area of electron diffraction (SAED) pattern obtained from CeO₂ film. It showed ring pattern characteristics revealing polycrystalline cubic crystal structure. From the SAED pattern, the lattice vector (G) and then d spacings were calculated. The SAED pattern was indexed and it matches with polycrystalline and cubic structure of CeO₂ as shown in Fig. 1. Fig. 3c shows the lattice image of nanocrystalline cubic CeO₂ of (200) orientation. This result confirmed the polycrystalline character of CeO₂ films with (200) preferred orientation and in good agreement with XRD results [26].

3.1.4. Surface morphology analysis

Atomic force microscopy was used to investigate the surface morphology and surface roughness of the films deposited on Si (100). Fig. 4A shows the surface morphology of the films

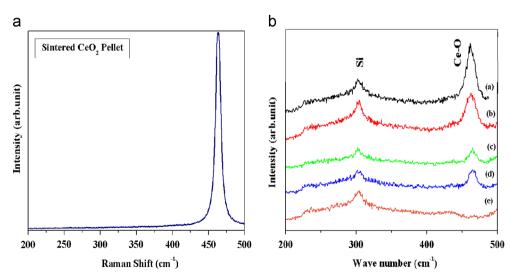


Fig. 2. (a) Raman spectrum of sintered ceria pellet; (b) Raman spectra of ceria thin films prepared at (a) 2.0×10^{-5} mbar; (b) 2×10^{-4} mbar; (c) 3.0×10^{-3} mbar; (d) 3.0×10^{-2} mbar; (e) 3.0×10^{-1} mbar.

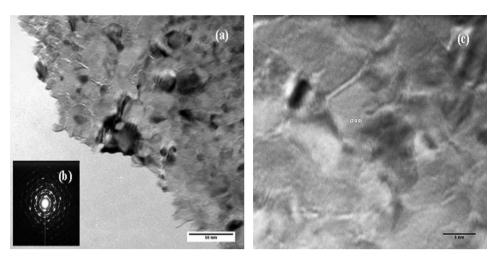


Fig. 3. (a) Low magnification bright field image of ceria film deposited on NaCl substrate at 73 K. (b) Selected Area of Electron Diffraction (SAED) pattern of ceria film. (c) Lattice image from ceria film showing (200) orientation.

prepared at an oxygen partial pressure range 2×10^{-5} 3×10^{-1} mbar. The surface morphology of the films change with thickness, substrate temperature and processing parameters, due to the change of the growth mode [27]. The films showed columnar growth, when the films prepared in the oxygen pressure range 2×10^{-5} – 3×10^{-3} mbar. At 3×10^{-2} mbar, the film showed the mixture of columnar and island growth. On increasing the oxygen pressure the columnar growth was observed to be suppressed. Therefore, the oxygen pressure has much influence on the energy of ablated species reaching the substrate and subsequently, the mobility of adatoms on the substrate surface. The other reported results showed that the average energy of ablated particles is almost constant in the pressure range 10^{-5} – 10^{-2} mbar, and it decreases considerably in the oxygen pressure $> 10^{-2}$ mbar, due to scattering and collision. Thus, increasing the oxygen pressure more than certain level, lead to deterioration in the crystallinity of films. The surface roughness of the films could be related to many parameters such as substrate temperature, ambient gas, deposition rate, repetition rate, energy density etc. The surface roughness was quantitatively measured by the root-mean-square (RMS) roughness. The surface roughness of the films increased (Fig. 4B) from 0.8 nm to 4.6 nm with increasing oxygen partial pressure from 2×10^{-5} mbar to 3×10^{-1} mbar, owing to the increased crystallinity [28].

3.2. Optical characterization

3.2.1. Photoluminescence spectroscopy

Photoluminescence (PL) spectroscopy was employed to study the optical properties of CeO_2 films deposited on Si (100). CeO_2 films were excited by He–Cd of 325 nm wavelength. The photoluminescence of all the films were measured at room temperature. The PL of CeO_2 films deposited at different oxygen partial pressure is shown in Fig. 5. The PL spectra showed a broad and strong peak \sim 400 nm and a weak shoulder peak \sim 442 nm. The PL measurements indicated the oxygen vacancy mediated emission from CeO_2 [29]. The broad PL of the as-deposited films could be the result of defects

including oxygen vacancies in the crystal with electronic energy levels below the 4f band. These defects possibly act as recombination centers for electrons initially excited from the valence band to the 4f band of the oxide [29,30]. The peak intensity of PL spectra of these samples decreased with increasing oxygen partial pressure, due to reduction in the concentration of O_2 related defects.

We noticed a broad PL peak at an energy value of 3.05 eV (~ 408 nm) for the films deposited in the oxygen pressure range 2×10^{-5} -3 $\times 10^{-3}$ mbar. It is found that the bandgap values increased from $3.05 \, \mathrm{eV}$ to $3.1 \, \mathrm{eV}$ as the oxygen pressure increased from 2×10^{-5} to $3 \times 10^{-1} \, \mathrm{mbar}$. There are defect energy levels between Ce4f and O2p level, so wide PL peaks could be observed. It was inferred that the PL of CeO2 is due to the transition by Ce4f \rightarrow O2p and defects level \rightarrow O2p [31,32]. These defects levels are located in the range of 1 eV around Ce4f band. The peak ~400 nm could be due to either 5d to 4f or 4f to valence band transitions [33] and coincides in energy with the 4f to valence band energy gap of 3.1 eV obtained from optical absorption measurements on CeO₂ reported in the literature [34– 36]. Giardini Guidoni et al. [36] deposited CeO₂ films on glass substrates at 400 °C by PLD and found that the energy bandgap of 3.1 eV. Ghanashyam Krishna et al. [37] revealed that the pure CeO₂ exhibited the bandgap of 3.3 eV at 873 K. Ozer [5] prepared the CeO2 films and found the optical bandgap value of 3.1 eV. Choudhury and Choudhury [38] prepared pure and doped CeO₂ nanoparticles and found the band to band excitation peak at 408 nm. Ansari [30] prepared the CeO₂ films and found the strong band \sim 378 nm in the photoluminescence spectra, correspond to the direct bandgap of 3.23 eV. These reported results are consistent with present results.

3.2.2. UV-visible spectroscopy studies

The refractive indices were calculated from the transmission spectra of CeO_2 films deposited on glass substrates at different oxygen partial pressure [28]. The refractive indices of the films decreased from 2.41 to 1.72 (at 550 nm wavelength) with increasing oxygen partial pressures and the values are given in Table 1. The variation of refractive indices is due to the

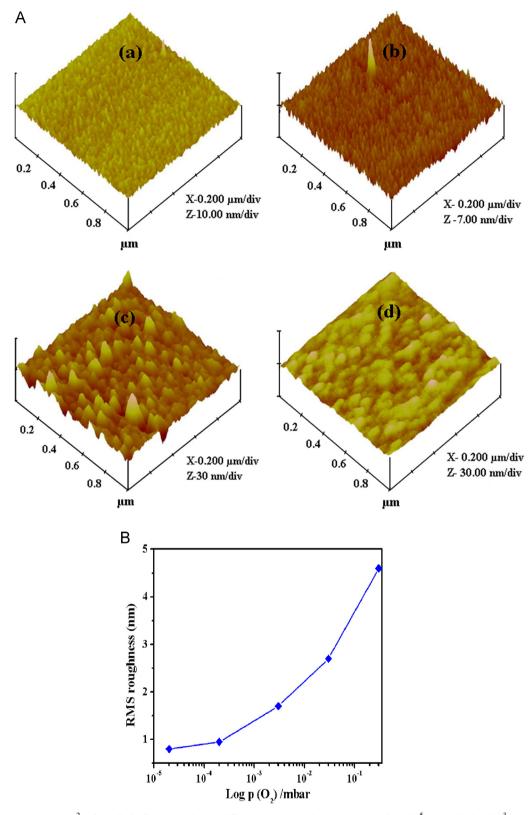


Fig. 4. (A) AFM images $(1 \times 1 \ \mu\text{m}^2)$ of the CeO₂ films deposited at different oxygen partial pressures: (a) 2×10^{-5} mbar; (b) 3×10^{-3} mbar; (c) 3×10^{-2} mbar; (d) 3×10^{-1} mbar. (B) RMS roughness of the films deposited as a function of oxygen partial pressure.

variation of packing density of the films. The packing density of the films depends on the energy of the ablated species. It is possible to correlate the oxygen partial pressure with the difference in the energy of the ablated species. At higher oxygen partial pressure, the energy of the ablated species is lower due to scattering and multiple collisions, causes films of lower thickness and packing density with rougher surfaces resulting lower refractive indices. Therefore, the higher oxygen

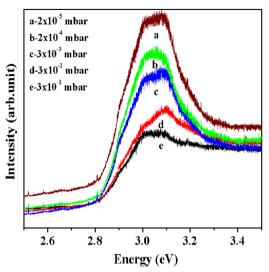


Fig. 5. PL spectra of ${\rm CeO_2}$ films deposited at various oxygen partial pressures and 673 K.

partial pressure causes lower packing density, whereas at lower oxygen partial pressure the energy of the species is higher, which leads to higher packing density of the films. Hence the refractive indices decreased with increasing oxygen partial pressure. The results from UV–visible spectrophotometer analysis showed that the optical properties of the films was influenced by the oxygen partial pressure [28,39,40]. The present investigation showed that the structure, preferred orientation, surface morphology and optical properties of the CeO₂ films could be controlled by a proper selection of the oxygen partial pressure, resulting high packing density and hence it can be used for optical applications.

4. Conclusions

The CeO_2 films were deposited on Si (100) and glass substrates with various oxygen partial pressures at a substrate temperature of 673 K by PLD. The XRD results showed that the films were polycrystalline cubic structure with (200) preferred orientation. The Raman peak $\sim 463~\rm cm^{-1}$ indicated the F_{2g} active mode of CeO_2 with a cubic structure. The HRTEM analysis confirmed the polycrystalline and cubic nature of CeO_2 films with (200) orientation. The AFM investigation indicated the change of growth mode of the films from columnar to island with increasing oxygen partial pressure range $2 \times 10^{-5} \text{--} 3 \times 10^{-1}$ mbar. The optical studies indicated that the optical bandgap values increased from 3.05 to 3.1 eV and the refractive indices decreased from 2.41 to 1.72 (at 550 nm wavelength) with increasing oxygen partial pressures.

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