

Short communication

In-situ diffuse phase transition at the Curie point of BaTiO₃ induced by amphoteric Ce³⁺/Ce⁴⁺ ions

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Abstract

Nominal (Ba_{1-x}Ce_x)(Ti_{0.95-x/4}Ce_{0.05})O₃ ceramics were prepared by a solid state reaction method. A complete solid solution called BC3TC5 was formed only at $x=0.03$, with a single-phase perovskite structure. The structural and dielectric properties of BC3TC5 were investigated by X-ray diffraction, scanning electron microscopy, Raman spectroscopy, electron paramagnetic resonance, and dielectric measurements. BC3TC5 exhibited a cubic structure and a non-uniform microstructure. Ti-vacancy defects exist in BC3TC5, which has a feature of mixed valence of Ba-site Ce³⁺ and Ti-site Ce⁴⁺. The amphoteric Ce³⁺/Ce⁴⁺ ions can induce an *in-situ* diffuse phase transition at the Curie point of BaTiO₃ ($T_C=125\text{ }^{\circ}\text{C}$). BC3TC5 exhibits a high- k behavior ($\epsilon'_m=9470$) and low dielectric loss ($\tan \delta < 0.025$).
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1. Introduction

As a candidate for capacitor applications, a Y5V-type Ba(Ti_{1-y}Ce_y)O₃ ceramic occurred at $y=0.2$ and showed a low rare-earth efficiency and dielectric permittivity ($\epsilon'_m \sim 2100$) [1]. In terms of La doping in Ba(Ti_{1-y}Ce_y)O₃, a high- k ‘Y5V’ ceramic (Ba_{1-x}La_x)(Ti_{1-y-x/4}Ce_y)O₃ ($x=0.03$, $y=0.05$) with room-temperature permittivity of $\epsilon'_{RT} > 10000$ can be achieved and the rare-earth efficiency can be raised significantly [2]. It is obvious that La doping at Ba sites plays an important role in a rapid peak shift toward low temperature and a rise in ϵ'_{RT} . (Ba_{1-x}La_x)(Ti_{1-y-x/4}Ce_y)O₃ can be considered as a solid solution of (Ba_{1-x}La_x)Ti_{1-x/4}O₃ and Ba(Ti_{1-y}Ce_y)O₃. The former contributes to the high- k behavior, while the latter can induce a strong diffuse phase transition (DPT) in the vicinity of room temperature. Since both (Ba_{1-x}La_x)Ti_{1-x/4}O₃ and (Ba_{1-x}Ce_x)Ti_{1-x/4}O₃ ($x \leq 0.05$) show a common feature of high- k first-order phase transition (FPT) [3–6], one holds

certain expectations about a new high- k solid solution (Ba_{1-x}Ce_x)(Ti_{1-y-x/4}Ce_y)O₃ by means of the formation of solid solution of (Ba_{1-x}Ce_x)Ti_{1-x/4}O₃ and Ba(Ti_{1-y}Ce_y)O₃, i.e., an amphoteric feature of Ce³⁺ at Ba sites (Ce³⁺_{Ba}) and Ce⁴⁺ at Ti sites (Ce⁴⁺_{Ti}) is reflected in the Ce doping behavior. It was reported that the solid solubility of Ce in BaTiO₃ is 8% for Ba-site Ce³⁺ incorporation, whereas it is 35% for Ti-site Ce⁴⁺ incorporation [1,7]. The preference difference in the site occupation will hinder the formation of a single-phase ceramic. It seems surprising that Ce ions play a role in both high- k and DPT behaviors. This will inevitably cause the separation of other phases such as CeO₂, BaCeO₃ [1], or Ce-rich Ba(Ti,Ce)O₃ phases [8] in the main perovskite phase.

On the other hand, we are very curious about the dielectric response caused by the amphoteric Ce³⁺_{Ba}–Ce⁴⁺_{Ti} ions in BaTiO₃ because the dielectric properties associated with the R³⁺_{Ba}–R⁴⁺_{Ti} (where R is any element) double substitution in BaTiO₃ has never been reported to date.

In this work, the nominal (Ba_{1-x}Ce_x)(Ti_{0.95-x/4}Ce_{0.05})O₃ ceramics were prepared to seek a critical balance of the mixed valence of Ce³⁺/Ce⁴⁺. Fortunately, a single-phase ceramic can be formed only at $x=0.03$. The present communication reports

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the structural and dielectric properties of this ceramic. It was found that an *in-situ* high-*k* diffuse phase transition occurred at the Curie point of BaTiO₃.

2. Experimental procedure

The starting materials were reagent-grade BaCO₃ (99.4%), TiO₂ (99.5%), and CeO₂ (99.9%). The ceramics were prepared according to the nominal formula (Ba_{1-x}Ce_x)(Ti_{0.95-x/4}Ce_{0.05})O₃ ($x=0.01$ – 0.05) using a mixed oxide method described elsewhere [2]. The final sintering conditions were chosen as 1450 °C for 24 h. Powder X-ray diffraction (XRD) measurements were made using a DX-2700 X-ray diffractometer (Dandong Haoyuan Inc.). Lattice parameters and unit cell volume were calculated by MS Modeling (Accelrys Inc.) using Cu K α 1 radiation ($\lambda=1.540562$ Å).

For single-phase (Ba_{1-x}Ce_x)(Ti_{0.95-x/4}Ce_{0.05})O₃ ($x=0.03$) ceramic (abbreviated BC3TC5), scanning electric microscope (SEM) images were obtained using a JSM-6490 SEM (JEOL) operated at 25 kV. Temperature dependences of the dielectric permittivity and loss were measured at 1 kHz using an RCL meter (Fluke PM6306). Raman spectra of ceramic powders were measured using a LabRAM XploRA Raman spectrometer (Horiba Jobin Yvon), with a 532 nm laser. A Linkam-600 heating and cooling stage was equipped for temperature-dependent Raman measurements. The electron paramagnetic resonance (EPR) spectrum was measured at an X band frequency of 9.84 GHz and at room temperature in the sweep range of 1500–5500 G using an EMX Plus spectrometer (Bruker). The gyromagnetic value (*g*) was calculated by the relationship $h\nu_0=g\beta H$.

3. Results and discussion

XRD patterns of the nominal (Ba_{1-x}Ce_x)(Ti_{0.95-x/4}Ce_{0.05})O₃ ceramics are shown in Fig. 1(a). The main perovskite phase for all the samples has been formed, showing a pseudo-cubic structure. A small amount of CeO₂ and BaCeO₃ was separated out from the main perovskite phase when $x\leq 0.02$, as clearly seen in Fig. 1(b) and (c). The Ce-rich perovskite phase of Ba

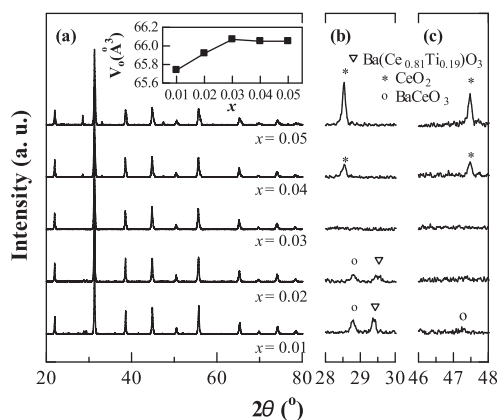


Fig. 1. XRD patterns of nominal (Ba_{1-x}Ce_x)(Ti_{0.95-x/4}Ce_{0.05})O₃ ceramics (a). The inset depicts variation in unit cell volume (V_0), as a function of x . (b,c) Enlarged XRD peaks in the vicinity of 29° and 47°, respectively for observation of the formation of a single-phase ceramic (BC3TC5).

(Ti_{0.81}Ce_{0.19})O₃ [8] appeared when $x\geq 0.04$. A complete solid solution called BC3TC5 was found to be formed at $x=0.03$. With an increase in x , the unit-cell volume (V_0) of the main perovskite phase increases linearly up to $x=0.03$ (Fig. 1 inset), indicating a consecutive incorporation of Ce⁴⁺ ions into Ti sites on the basis of an ionic size comparison (12-CN Ba²⁺: 1.61 Å, Ce³⁺: 1.34 Å; 6-CN Ti⁴⁺: 0.605 Å, Ce⁴⁺: 0.87 Å [9]). The maximum value of V_0 appeared at $x=0.03$. Subsequently, a slower decrease in V_0 for $x\geq 0.03$ implies a decrease in the solubility of Ce⁴⁺ at Ti sites. Hence, the formula of (Ba_{1-x}Ce_x)(Ti_{0.95-x/4}Ce_{0.05})O₃ is tenable just for $x=0.03$. Makovec et al. suggested that cerium is incorporated into the BaTiO₃ lattice at Ba sites as Ce³⁺ and at Ti sites as Ce⁴⁺ in the ratio of 4:3 when pure CeO₂ is added to BaTiO₃ [7]. The ratio of our single-phase sample with $x=0.03$ is 3:5. The difference between the two results arises from the differences in the Ba/Ti ratio and in the sintering temperature.

The V_0 value of BC3TC5 (66.08 Å³) is less than that of Ba(Ti_{0.95}Ce_{0.05})O₃ (66.37 Å³) [2], but greater than that of the cubic BaTiO₃ (65.50 Å³) (JCPDS Cards no. 31-174). This fact gives an evidence for the existence of the mixed valence of Ba-site Ce³⁺ and Ti-site Ce⁴⁺ in BC3TC5, i.e., the formation of Ce³⁺_{Ba}–Ce⁴⁺_{Ti} complexes.

Temperature-dependent Raman spectra of BC3TC5 are shown in Fig. 2. The two broad bands at 260 and 525 cm⁻¹, which persisted till 180 °C, are attributed to intrinsic disorder in the cubic phase of BC3TC5 [10]. The rapid attenuation of the 720 cm⁻¹ band at 180 °C implies an increase in cubicity. As with the Raman spectrum observed in (Ba_{1-x}La_x)(Ti_{0.95-x/4}Ce_{0.05})O₃ ($x=0.03$) ceramics (BL3TC5) [10], a weak 824 cm⁻¹ band, which was generally caused by trivalent rare-earth ions (La³⁺, Nd³⁺, Eu³⁺) at Ba²⁺ sites in the BaTiO₃ lattice [11–13], also gives evidence for Ce³⁺ entering Ba sites in BC3TC5. However, the 824 cm⁻¹ band of BC3TC5 is much weaker

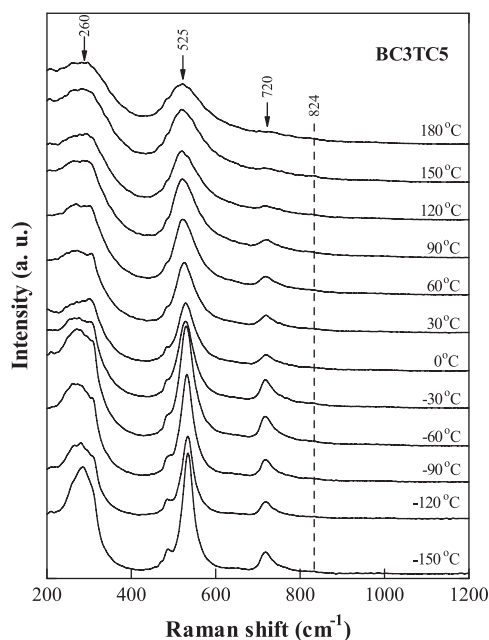


Fig. 2. Temperature-dependent Raman spectra of BC3TC5.

compared to BL3TC5 [10], which suggests that the Raman activity induced by Ba-site Ce^{3+} due to the formation of $\text{Ce}_{\text{Ba}}^{3+}-\text{Ce}_{\text{Ti}}^{4+}$ complexes is dramatically reduced.

EPR spectrum of BC3TC5 is shown in Fig. 3. A signal at $g=2.003$ was assigned to ionized Ti-vacancy defects (V_{Ti}) [10,14], which was caused by the substitution of Ce^{3+} at Ba^{2+} sites, in good agreement with the Raman scattering observations.

SEM image of BC3TC5 is shown in Fig. 4. BC3TC5 exhibited a non-homogeneous but denser microstructure consisting of coarse grains (3–6 μm) and fine grains ($\sim 1 \mu\text{m}$). The average grain size is 2.7 μm . The morphology of the other four samples with $x \neq 0.03$ is the same as BC3TC5 (not presented here). The secondary phase that grew between three or four coarse grains boundaries can be observed when $x \geq 0.04$.

Temperature dependence of the dielectric permittivity (ϵ') and loss ($\tan \delta$) for BC3TC5 is shown in Fig. 5. The $\epsilon'-T$ curves of BL3TC5 and BaTiO_3 are depicted in this figure for comparison [2]. BaTiO_3 showed a first-order phase transition and a sharp-peak feature at its Curie point ($T_C \approx 125^\circ\text{C}$). For BL3TC5, the so-called Curie temperature (T_C) at which the maximum dielectric permittivity (ϵ'_m) occurred decreased rapidly at a dramatic rate of $-30^\circ\text{C}/\text{at}\% \text{La}$ [2]; therefore the dielectric peak of BL3TC5 could approach room temperature ($T_C = 38^\circ\text{C}$) when x was increased to 0.03. Similar to BL3TC5, BC3TC5 also exhibited a strong high- k diffuse phase transition (DPT) behavior ($\epsilon'_m = 9470$), which demonstrates that Ba-site Ce^{3+} ions in BC3TC5 are responsible for raising

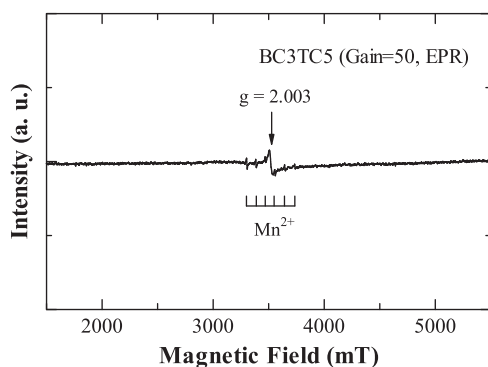


Fig. 3. EPR spectrum of BC3TC5.

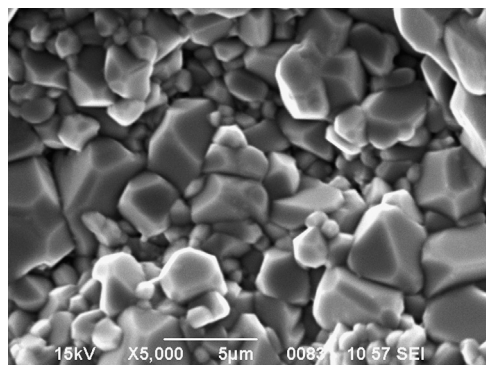


Fig. 4. SEM morphology of BC3TC5.

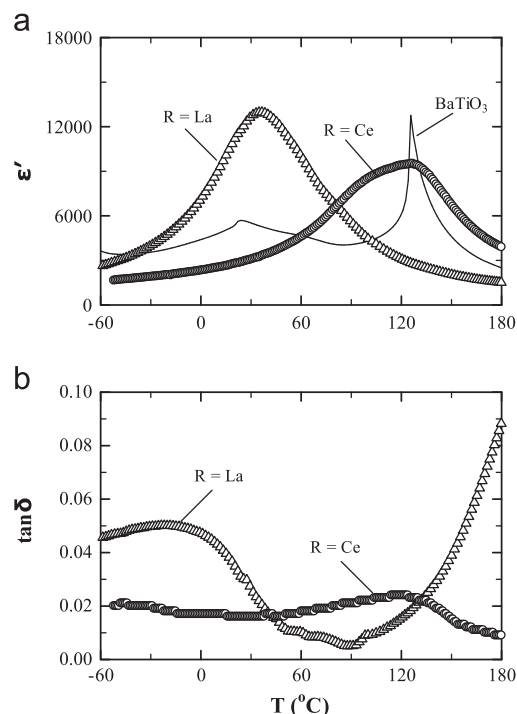


Fig. 5. Temperature dependence of (a) the dielectric permittivity (ϵ') and (b) the dielectric loss ($\tan \delta$) for BC3TC5 and BL3TC5. The data of BL3TC5 and BaTiO_3 are from Ref. [2].

the dielectric permittivity and Ti-site Ce^{4+} ions play a predominant role in the DPT behavior [2].

However, the T_C of BC3TC5 occurred at 125°C . This fact indicates that the amphoteric $\text{Ce}^{3+}/\text{Ce}^{4+}$ ions make no contribution to a peak shift toward low temperature, and only induce an *in-situ* diffuse phase transition at the Curie point of BaTiO_3 . The other four samples with $x \neq 0.03$ also exhibited high- k DPT behavior ($\epsilon'_m > 9000$) and their T_C occurred at $123\text{--}127^\circ\text{C}$ (not presented here). Thus, *in-situ* DPT is the common property of this series of ceramics.

At room temperature, BC3TC5 exhibits a cubic structure, which is not in contradiction with T_C observed at 125°C . Because of the nature of DPT, BC3TC5 can retain an average cubic structure at lower temperatures far from T_C , as observed for $(\text{Ba}_{1-x}\text{La}_x)(\text{Ti}_{0.95-x/4}\text{Ce}_{0.05})\text{O}_3$ ($x=0.02$) ($T_C=64^\circ\text{C}$) [2]. As far as phase structure is concerned, BC3TC5 with DPT is different from BaTiO_3 with first-order phase transition (FPT). The latter shows a very sharp dielectric peak at its T_C , accompanied by a clear cubic–tetragonal phase transition, whereas the dielectric peak of the former (BC3TC5) becomes vague. In the case of DPT, the so-called T_C , strictly speaking, cannot represent the cubic–tetragonal phase transition point, i.e., the Curie temperature. At present, T_C is often represented by the dielectric-peak temperature (T_m).

A type of ions that can split into the two sites in BaTiO_3 occurs in very few cases. This type of ions is usually trivalent. They form a self-compensation (donor–acceptor) mode of $\text{R}_{\text{Ba}}^{3+}-\text{R}_{\text{Ti}}^{3+}$ in $(\text{Ba}_{1-x}\text{R}_x)(\text{Ti}_{1-x}\text{R}_x)\text{O}_3$ such as trivalent rare-earth ions $\text{R}=\text{Nd}$ [15], Eu [13], and Er [16,17]. All of this type of ions can induce strong DPT, accompanied by a peak shift toward low temperature

[13,15,16]. The dual character of $\text{Ce}^{3+}/\text{Ce}^{4+}$ in BaTiO_3 is the sole case of the $\text{R}_{\text{Ba}}^{3+}-\text{R}_{\text{Ti}}^{4+}$ (R is any element) double substitution to date. The mechanism of *in-situ* DPT of BC3TC5 is not known.

BC3TC5 exhibited lower dielectric loss ($\tan \delta < 0.025$). A stronger Mn^{2+} ($3d^5$) sextet signal was observed (Fig. 3), which indicates that a large number of Mn^{4+} or Mn^{3+} impurities in BC3TC5 were reduced into Mn^{2+} . There is almost no Mn^{2+} signal in the EPR spectrum of BL3TC5 [10]. Thus, the lower dielectric loss of BC3TC5 is attributed to the role of trapping electrons by Mn impurities.

4. Conclusions

The sole single-phase solid solution (BC3TC5) with $x=0.03$ can be prepared using a solid state reaction method according to the formula $(\text{Ba}_{1-x}\text{Ce}_x)(\text{Ti}_{0.95-x/4}\text{Ce}_{0.05})\text{O}_3$. BC3TC5 exhibits a cubic structure and a non-uniform microstructure. BC3TC5 with Ti-vacancy defects has a feature of mixed valence of Ba-site Ce^{3+} and Ti-site Ce^{4+} in the ratio of 3:5. A strong diffuse phase transition due to the amphoteric $\text{Ce}^{3+}/\text{Ce}^{4+}$ occurs at the Curie point of BaTiO_3 ($T_C=125^\circ\text{C}$). BC3TC5 shows a high- k behavior ($\epsilon'_m=9470$) and low dielectric loss ($\tan \delta < 0.025$). The ability of trapping electrons by Mn impurities is responsible for the lower $\tan \delta$ of BC3TC5.

Acknowledgments

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