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Effects of nitrogen pressure on microstructures and properties of Sr₂Fe_{10/9}Mo_{8/9}O₆ thin films

Jie Xu, Wei-Jing Ji, Jing-Feng Wang, Zheng-Bin Gu*, Shan-Tao Zhang*

National Laboratory of Solid State Microstructure and Department of Materials Science and Engineering, Nanjing University, Nanjing 210093, China

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Abstract

High quality double perovskite $Sr_2Fe_{10/9}Mo_{8/9}O_6$ thin films have been prepared on (111)-SrTiO₃ substrates under different N_2 partial pressures (0.1–20 Pa). The films are (111)-oriented with flat surfaces and sharp film/substrate interfaces. The epitaxy and local Fe/Mo ordering are revealed by high resolution transmission electron microscopy and selected area electron diffraction. It is established that the N_2 pressure has significant effects on structures and properties: With increasing N_2 pressure, the film thickness decreases and the strain increases, while the magnetization decreases monotonously. Interestingly, all the films show semiconductor-like behavior in the whole measuring temperature range but the resistivity increases with increasing N_2 pressure. The observed results are attributed to that N_2 pressure can affect the strain and the concentration of Fe–O–Mo, both play important roles in magnetic interaction, and thus affect the magnetic and transport behaviors.

Keywords: A. Films; Double perovskite; Ferromagnetism; Microstructure

1. Introduction

Ferromagnetic double perovskites have been extensively studied due to their fascinating fundamental physics and their potential for applications in devices such as spintronics, magnetic sensors, etc. [1,2]. These oxides can show high Curie temperature (T_c), good magnetic and transport properties even at room temperature (RT). As a typical example for ferromagnetic double perovskite, half-metal Sr_2FeMoO_6 (SFMO) has T_c of 420 K and magnetroresistance (MR) of -10% at RT [2,3]. Theoretically, SFMO has a net magnetization of 4 μ_B /formula unit (f.u.), resulting from the antiferromagnetical coupling between Fe^{3+} (3d⁵, S=5/2) and Mo^{5+} (4d¹, S=1/2) [2,3].

Up to now, pure SFMO and A- and/or B-site substituted SFMO have been intensively studied in both ceramic and thin film forms [3–9]. However, one of the drawbacks for SFMO-based double perovskite is that reducing atmosphere is

E-mail addresses: zbgu@nju.edu.cn (Z.-B. Gu), stzhang@nju.edu.cn (S.-T. Zhang).

necessary for preparing single-phase, high quality samples. This drawback leads to both safety and cost considerations. Therefore, an alternative double perovskite composition, which can show ferromagnetism above RT and can be synthesized without reducing atmosphere, is desirable.

Recently, $Sr_2Fe_{1+x}Mo_{1-x}O_6$ (x=1/9), abbreviated as SFMO-1/9, is carefully designed and synthesized under relatively safe and cheap N_2 [10]. This composition is found to show ferromagnetism with T_c of 390 K and MR of -10% (10 K, 2 T external magnetic field). Additionally, different from the most reported, not well-designed off-stoichiometric $Sr_2Fe_{1+x}Mo_{1-x}O_6$, this composition can theoretically avoid intrinsic, off-stoichiometry-induced oxygen defects since it is actually the mixture of stoichiometric SFMO and stoichiometric $Sr_3Fe_2MoO_9$ (i.e., $Sr_2Fe_{4/3}Mo_{2/3}O_6$) with the ratio of 2:1.

Accordingly, SFMO-1/9 might be a promising room temperature ferromagnetic double perovskite without safety and cost considerations. Based on the work on ceramics [10], SFMO-1/9 thin films with different thicknesses have been reported [11]; however, it is noted that the atmosphere pressure during film fabricating is also one of the key factors determining the mechanical strength, microstructures, and thus macroscopic

^{*}Corresponding authors.

properties [12–14]. That is to say, detailed works on the effects of depositing atmosphere pressure on structures and properties of SFMO–1/9 films are necessary to further optimize the properties. For this goal, in this paper, SFMO–1/9 films have been prepared in N_2 stream with different pressures of 0.1–20 Pa, the effects of N_2 pressure on structures and properties are discussed in detail.

2. Experimental

The SFMO – 1/9 target for pulsed laser deposition (PLD) was synthesized by solid state reaction with starting materials of SrCO₃, Fe₂O₃, and MoO₃. Detailed information about target preparation can be found elsewhere [9]. The SFMO – 1/9 thin films were fabricated on (111) SrTiO₃ (STO) substrates at 750 °C for 30min by using the 355 nm output of a Brilliant Nd: YAG laser. During deposition, the laser repetition rate is 10/3 Hz, the pulse energy is 200 mJ, and the partial N₂ pressure varied from 0.1 to 20 Pa. After deposition, the films were *in situ* annealed in chamber for 10min.

The structures of the films were investigated by X-ray diffraction (XRD, Rigaku Ultima III) and transmission electron microscopy (TEM, FEI Tecnai F20). The cross-sectional TEM samples were prepared by mechanical thinning and ion milling (Gatan PIPS) with 3.2 keV – 4.4 keV Ar⁺ ions for few hours. The magnetic and transport data were collected by using a superconductor quantum interference device (SQUID, Quantum Design, MPMS XL-7) and a physical property measurement system (PPMS, Quantum Design, 2001NUGC).

3. Results and discussion

Fig. 1 shows the X-ray $\theta - 2\theta$ -scans of the films deposited at 0.1, 10, 15 and 20 Pa N₂. As can be seen, only (222) peaks from the films and (111) peaks from the substrates are observed, indicating that the films are well (111)-oriented. Additionally, under 0.1 Pa and 10 Pa N₂ pressures, the diffraction peaks from the films are observed clearly locating

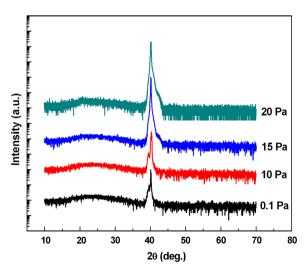


Fig. 1. X-ray diffraction patterns for the films deposited under different N_2 pressures.

left to the peaks from the substrates. The pseudocubic lattice constants are estimated to be 7.9121 and 7.8746 Å, and accordingly, the strains are estimated to be from 1.31% to 0.82%, respectively, by noting the cubic lattice constant of STO is 3.905 Å. However, when N_2 pressure is increased to 15 and 20 Pa, the peak from the films tends to combine with that from the substrates. This observation suggests that the film strain increases with increasing N_2 pressure, consistent with other report that partial pressure can lead to significant mechanical strength changing in thin films [14].

The microstructures of the films deposited at 0.1 and 15 Pa N₂ are typically investigated by TEM. The low magnification cross-sectional TEM micrographs are shown in Fig. 2(a) and (b). In general, both films are homogeneous and have very smooth surfaces and sharp film/substrate interfaces. Columnlike microstructure indicates the three-dimensional growth mode, which is typical for films deposited by the PLD method. The average thickness of the films deposited at 0.1 Pa and 15 Pa N₂ is slightly different, with the values of 48 nm and 43 nm, respectively, therefore, higher the deposition pressure, smaller is the film thickness. This result is reasonable since the laser ablated compositions will have interaction with the deposition atmosphere during transporting to the substrate, higher deposition atmosphere pressure means stronger and more interactions, and thus less composition can reach the substrate to form films. Based on this TEM result, thickness of the films deposited at 10 Pa and 20 Pa N₂ are reasonably estimated to be close to 43 nm and the film thickness decreases monotonously with increasing N₂ pressure.

High resolution TEM (HRTEM) images close to the film/substrate interface for the films deposited at 0.1 and 15 Pa N_2 are taken along [011] zone-axis, as shown in Fig. 3(a) and (b), respectively. The arrows indicate the film/substrate interfaces. As can be seen, the interfaces are very sharp and the atomic arrangements across the interfaces are coherent, indicating high quality and epitaxial growth of the films.

The selected area electron diffraction (SAED) images of the films deposited at 0.1 and 10 Pa N_2 are taken from [011] zone-axis, indexed based on pseudo-cubic setting and shown in Fig. 4(a) and (b), respectively. Superlattice diffraction spots are observable, as indicated by the arrows. These superstructure reflection spots are $1/2\{h+1/2,k+1/2,l+1/2\}$, and arise from

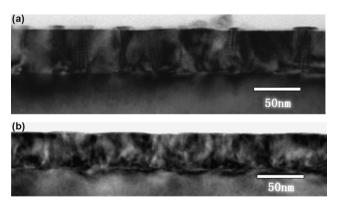


Fig. 2. The low-magnification cross-sectional TEM images of SFMO-1/9 films deposited at (a) 0.1 Pa, (b) 15 Pa N_2 pressure.

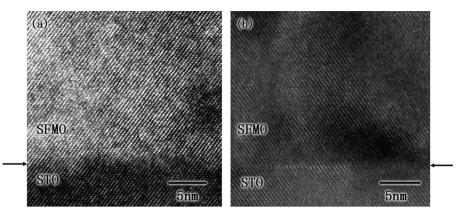


Fig. 3. HRTEM images of the interface, taken from [011] zone-axis, for the films deposited at (a) 0.1 Pa, (b) 15 Pa N_2 pressure.

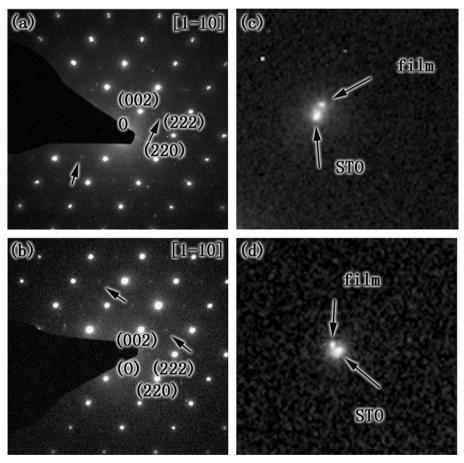


Fig. 4. SAED for the films deposited at (a) 0.1 Pa, (b) 15 Pa N_2 pressure. The enlarged diffraction spots from films and substrates show films deposited at 15 Pa N_2 (c) have higher strain than that deposited at 0.1 Pa N_2 .

the local Fe/Mo ordering which makes the unit cell twice in real space [7,15,16]. Additionally, from Fig. 4(c) and (d), it is found that the films deposited at 15 Pa have closer reflection spots from substrates and films than that deposited at 0.1 Pa N_2 , this may suggest that the films deposited under higher N_2 pressure have higher strain, consistent with the XRD results.

Magnetization-magnetic field (M–H) loops for the films deposited at 0.1, 10, 15 and 20 Pa N_2 are measured at 10 K and RT with the field of 2 T, the loops are illustrated in Fig. 5(a) and (b), respectively. In general, the films show well-saturated

M–H ferromagnetic hysteresis loops. At both temperatures, the saturation magnetization ($M_{\rm s}$) decreases slightly with increasing N₂ pressure. For example, at 10 K, the values of $M_{\rm s}$ are 0.38, 0.35, 0.32, and 0.27 μ_B/f.u. for the films deposited at 0.1, 10, 15, and 20 Pa N₂, respectively. As can be seen, for the films deposited under 0.1 and 10 Pa N₂, the strain decreases from 1.31% to 0.82%, and the magnetization at 10 K decreases by 7.9% from 0.38 to 0.35 μ_B/f.u. Such pressure-dependent magnetization may be attributed to the following two factors. On the one hand, the strain increases with increasing N₂

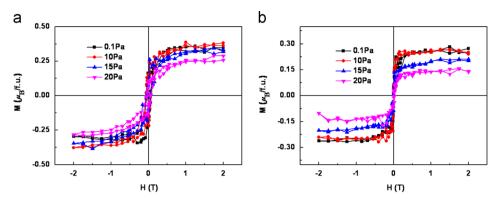


Fig. 5. M-H loops of the films at (a) 10 K and (b) 300 K.

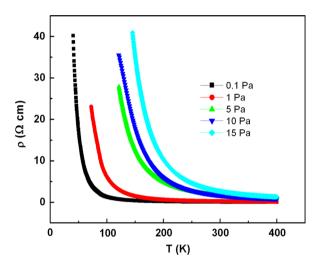


Fig. 6. Temperature dependence of resistivity of the films deposited under different N_2 pressures.

pressure, as shown in Figs. 1 and 4, the changed strain will affect the bond length and the bond angle of Fe–O–Mo. And on the other hand, some N may occupy the O sites, so N can enter into lattice, higher N₂ pressure means more substituted O by N, which might weakens the magnetic interaction by decreasing the concentration of Fe–O–Mo. These two factors can have considerable effects on magnetic interaction, and thus the macroscopic magnetic and transport behaviors [4,15,16]. To deeply understand how the strain affects bond length, bond angle, and the substituting of O by N, detailed investigating and analyzing are necessary.

From Fig. 5, it is found that the experimental magnetization value is far lower than the theoretical value (4.0 $\mu_B/f.u.$ for pure SFMO). As well known, the film thickness and the antisite disorder (ASD) play an important role in determining macroscopic magnetization value. In general thicker films and higher ASD content will lead to smaller magnetization [11]. The facts that out SFMO-1/9 should have intrinsic ASD because the Fe/Mo ratio deviates from 1:1 and the film thickness which can affect magnetic interaction as described above, should be responsible for the observed significantly low magnetization. Actually, 11 nm SFMO-1/9 films have a large magnetization of 2.3 $\mu_B/f.u.$ [11].

The films deposited under 20 Pa N₂ have too large resistance to be measured. The resistances of the films deposited at 0.1, 1, 10, 15 Pa N₂ are measured and plotted in Fig. 6 as the functions of temperature. In general, all the films show semiconductor-like behavior in the whole measuring temperature range. However, the resistivity increases with increasing N₂ pressure. At 150 K, the resistivity values are 0.5, 1.7, 12.5, 17.6, and 36.1 Ω cm for the films deposited at 0.1, 1, 5, 10, and 15 Pa N₂, respectively. Actually, this observation may be attributed to same reasons responsible for magnetization as discussed above. It is noted for the films deposited under 0.1 Pa and 10 Pa, the strain decreases from 1.31% to 0.82%; whereas the resistivity increases by 35 times from 0.5 to 17.6 Ω cm, implying the significant effects of strain on transport property. Further detailed works quantifying the relationship between mechanical strength, bond length/ bond angle, and magnetic/transport properties are interesting and may be helpful for further understanding the physics background.

4. Conclusions

In summary, SFMO-1/9 thin films have been prepared on SrTiO₃ (111) substrates under different N₂ pressures. The effects of N₂ partial pressure on microstructure, magnetic and transport properties have been established and discussed. It is experimentally found that with increasing N₂ pressure, the strain increases, the magnetization decreases, and the resistivity increase. These observed results are attributed to the fact that higher N₂ pressure leads to larger strain and more substituting of O by N, since the former can change the bond length and bond angle of Fe–O–Mo, while the later can decrease concentration of Fe–O–Mo, both have effect on magnetic interaction. As the results, macroscopic magnetic and transport behaviors show significant N₂ pressure dependence. Our results may be helpful for optimizing properties of ferromagnetic double perovskite thin films.

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