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High carrier density CeO₂ dielectrics—implications for MOS devices

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Abstract

Nanocrystalline CeO_2 films of 75 nm thickness were deposited on n-type (100) silicon substrates using pulsed laser deposition (PLD) technique to form a gate dielectric in the $Pt/Si/CeO_2/Pt$ MOS capacitor. XRD and AFM were used for characterization of the film crystal structure and grain size. Electrical properties of the MOS structure were examined by capacitance-voltage (C-V) and impedance spectroscopy measurements. Based on the impedance measurements and reported electron mobilities, we derive a rather high carrier density of the order of 6×10^{17} cm⁻³. Nevertheless, their overall high resistance enables the films to serve as dielectrics. In contrast to conventional MOS capacitors, we find an additional capacitive contribution under accumulation we attribute to the electron depletion in the CeO_2 film. We present a model consistent with these results.

Keywords: Capacitor; CeO2; Electrical properties; Films; Microstructure

1. Introduction

The continued drive towards ever smaller submicron lateral dimensions in MOSFET technology has pushed the SiO₂ gate oxide thickness to the order of nanometers leading, in turn, to higher leakage currents and processing difficulties. Replacement of the SiO₂ films by thicker films with higher dielectric constants compared with SiO₂ ($\varepsilon_r \approx 3.5$) promises equivalent capacitances with reduced leakage and threat of electrical breakdown. 1 Cerium dioxide (CeO₂) appears to be a particularly attractive candidate, given its high dielectric constant ($\varepsilon_r \approx 25$) and its compatibility with Si. In the solid state ionics arena, there is much interest in reducing the operating temperatures of solid state oxide fuel cells (SOFC) by shifting from bulk to thin film electrolytes (stabilized ZrO₂ or acceptor doped CeO₂). The lattice constant mismatch between diamond-structured silicon and fluorite-structured CeO2 is only about 0.35%, thus enabling the fabrication and deposition of stress free, highly oriented or epitaxial thin film structures. Furthermore, CeO₂ forms no undesirable reaction products, such as silicides, when in contact with silicon

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at elevated temperatures.^{2–4} Lastly, CeO₂ thin films are effective as buffer layers and diffusion barriers between the Si substrate and many complex and chemically reactive compounds, such as Pb(Zr,Ti)O₃ and PbTiO₃.^{5–7} In this paper we describe the fabrication and characterization of the Pt/Si/CeO₂/Pt MOS device structure. The electrical properties of the CeO₂ thin films are deconvoluted from the overall response of the MOS device with the assistance of impedance spectroscopy. The analysis reveals an interesting new property of the MOS device based on the features of high carrier density CeO₂ dielectrics.

2. Experimental

CeO₂ thin films were deposited onto n-type silicon wafers by a pulsed KrF excimer laser (Lambda Physik LPX-300) operating at a wavelength of 248 nm. The laser pulse energy was 400 mJ, with the beam fluence set to 3.65 J/cm² with aid of focusing optics. In situ deposition was carried out in a vacuum chamber 1.5×10^{-3} Torr at the temperature of 600 °C. n-Type silicon substrates with resistivity of 2–4 Ω cm and (100) surface orientation were placed parallel to the target at a distance of 65 mm from the target surface. Before deposition, the silicon substrates were etched with 5% HF in water and rinsed with deionized water. Both back

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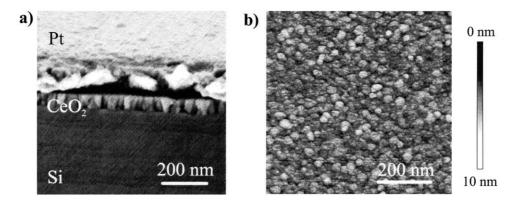


Fig. 1. (a) FESEM cross section micrograph of a Pt/Si(100)/CeO₂/Pt structure and (b) AFM micrograph of CeO₂ thin-film surface. Columnar structure of the CeO₂ film in (a) is clearly evident.

and top platinum electrodes with thickness of 150 nm, and area of 130×10^{-6} and 2.405×10^{-6} m², respectively, were deposited by room temperature DC sputtering. Back electrode entirely covered the backside of the sample.

The CeO₂ target preparation included cold pressing pure CeO₂ powder at a pressure of 20 MPa, vacuum sealing the resultant pellet, followed by isostatic pressing at 290 MPa. After sintering at 1425 °C for 10 h in air, the target had a density of 6.5×10^3 kg/m³. X-ray diffraction θ –2 θ measurements (Rigaku RU300) were performed on both CeO₂ films and target. A tapping mode AFM (Nanoscope IIIa) and SEM (Jeol JSM-6400) were used for the evaluation of thin film surface quality. In order to study the electrical properties of the MOS structure, an impedance analyzer with frequency sweep from 0.1 Hz to 1 MHz (Solartron 1260) was used for both the complex impedance spectroscopy and C–V measurements.

3. Results and discussion

SEM and AFM micrographs of the Si(100)/CeO₂/Pt multilayer cross section and the CeO₂ thin-film surface are shown in Fig. 1(a) and (b), respectively. Analysis of

X-ray diffraction measurements show the films to be mainly (111) oriented. Some contributions from (311) planes and reduced peak intensities suggest a reduced level of order when compared to epitaxial (111) oriented films.⁸ In Fig. 1(a), a columnar grain structure of the CeO₂ films is clearly revealed. AFM micrograph of the sample presented in Fig. 1(a) is shown in Fig. 1(b). Analysis of this micrograph revealed a rms roughness value of $R_q \approx 0.63$ nm. Also, from larger scale AFM micrographs, it was possible to estimate the grain size to be ~ 40 to 50 nm and to conclude that no large particulates were present, pointing to the high quality target used in the deposition process.

In order to correctly interpret the electrical response of the multilayer MOS device, one has to carefully consider the responses arising from the different regions of the device. For example, the potentially blocking nature of the Pt/Si and CeO₂/Pt contacts (work function differences: $\Phi_{\text{Pt/Si}} = 1.05 \text{ eV}$ and $\Phi_{\text{CeO}_2/\text{Pt}} = 0.6 \text{ eV}$), must to be taken into consideration as extra voltage-dependent capacitors connected in series with the dielectric film capacitance C_{ox} . In similar manner, one has to consider the space charge barrier layer resistance and capacitance at the Si/CeO₂ interface when interpreting the impedance spectrum. In Fig. 2(a) and (b), a

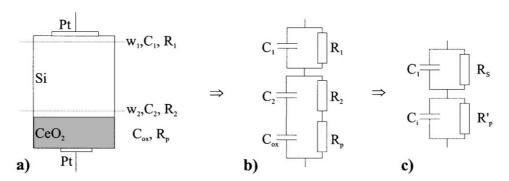


Fig. 2. Schematic drawing of the MOS device describing (a) the different space charge capacitances and resistances C_{1} , R_{1} and C_{2} , R_{2} , and the oxide capacitance and resistance C_{ox} , R_{p} together with corresponding equivalent circuit in (b). In (c) the equivalent circuit is rearranged to accent the responses of the capacitances C_{2} and C_{ox} typically studied in C–V measurements.

schematic structure of the MOS device and its equivalent circuit are sketched, respectively. In order to emphasize the important contributions for the C–V analyses, the equivalent circuit was rearranged according the Fig. 2(c) leading to the following equation for the impedance of the MOS structure:

$$Z = \left[\frac{R_{\rm s}}{1 + \omega^2 C_1^2 R_{\rm S}^2} + \frac{R_{\rm p}'}{1 + \omega^2 C_{\rm i}^2 R_{\rm p}^2} \right]$$

$$- j \left[\frac{\omega C_1 R_{\rm S}^2}{1 + \omega^2 C_1^2 R_{\rm S}^2} + \frac{\omega C_{\rm i} R_{\rm p}'^2}{1 + \omega^2 C_{\rm i}^2 R_{\rm p}'^2} \right]$$

$$= Z' + jZ''. \tag{1}$$

Since the dielectric constant of the silicon substrate is only 11.8 compared with that of the CeO₂ with $\varepsilon_r \sim 25$, and the thickness of the substrate is 500 µm compared with 75 nm of the film, respectively, the applied DC bias actually vanishes across the substrate space charge regions. Especially, in the state of accumulation, when the effect of C_2 in Fig. 2 disappears, C_1 can be estimated from the expression for the Schottky-junction capacitance.¹⁰ An approximate value for the oxide resistance $R'_{\rm p} = 12.6 \text{ k}\Omega$ was derived from the impedance spectroscopy measurements. The resistance $R_s = 100 \text{ M}\Omega$ was obtained as a function of the applied DC bias from impedance measurements carried out for the reference component without a CeO₂ thin film, i.e., Pt/Si/Pt. This resistance is not independent of DC bias, but represents typical values for back-to-back Schottky junction devices.

In Fig. 3, there are three C–V curves presented for a Pt/Si(100)/CeO₂/Pt MOS device measured at 100 Hz,

1 kHz, and 100 kHz, respectively. First of all, there are both depletion-inversion and accumulation plateaus shown in each C-V curve typical for MOS devices. At the frequencies of 100 Hz and 1 kHz, a value of 0.2 nF for the space charge capacitance C2 in depletion was measured. This is very close to the calculated value of capacitance for an n-type silicon substrate with $N_{\rm d} = 1.5 \times 10^{15}$ cm⁻³. An oxide layer capacitance $C_{\rm ox}$ = 6.5 nF for CeO₂ in accumulation was also measured, giving a dielectric constant of $\varepsilon_r = 23$. There is also a clear shift of the flat band voltage $V_{\rm FB}$ to higher gate bias voltages with increasing frequency pointing to a contribution from interface traps. This shift of the flat band voltage might be pronounced due to the high concentration of defects in nonstoichiometric CeO₂ films with nanosized grains impairing the properties of the interface when compared, for example, with SiO₂/Si interface. 10 Furthermore, there is a slight increase in oxide capacitance C_{ox} with decreasing frequency. This, we believe, is a consequence of the heterogeneous Maxwell-Wagner effect typical for ceramics having grain boundaries with different conductivity than bulk grains. This effect is more prominent at frequencies below 100 Hz not shown here. 11 In addition to the conventional C-V measurement response, there were also some unusual features observed at low frequencies. Above a certain positive gate bias, there is a strong increase in the value of capacitance $C_i = (1/C_2 + 1/C_{ox})^{-1} \approx C_{ox}$ in accumulation shown in Fig. 3. This kind of behavior was found also by Sakai et al. 12 for a similar component structure with ohmic Al contacts, and they suggested that this increase was caused by leakage in CeO₂. However, we found this phenomenon to be strongly frequency

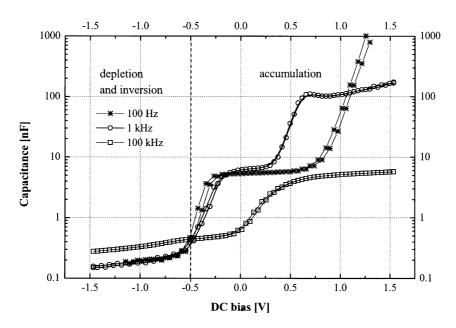


Fig. 3. Three C-V curves of a $Pt/Si(100)/CeO_2/Pt$ device calculated from impedance spectroscopy data measured at 100 Hz, 1 kHz, and 100 kHz. Notice the second accumulation capacitance plateau in the curve measured at 1 kHz generated due to the depletion of the space charge in CeO_2 .

dependent and thus probably related to space charge developed in the CeO₂ thin film, which might also cause leakage current under DC biasing conditions. At the frequency of 1 kHz, another plateau with capacitance of 100 nF was obtained. If one considers a biasing condition where the back-electrode Schottky barrier is totally reverse biased, the voltage drop across the depletion layer and the corresponding width of the depletion layer w_1 eventually saturates. When the gate potential is still increased, the back-electrode capacitance becomes meaningless in overall capacitance, and, thus, the second plateau of capacitance 100 nF in Fig. 3 must be located inside the CeO₂ film. The room temperature resistance of the CeO₂ film was $R_{\rm p}'=12.6~{\rm k}\Omega$ leading to a conductivity of $\sigma=2.467\times10^{-8}~{\rm S/cm}$. The electron density in the oxide can be estimated from knowledge of the published electron mobility. Using the measured activation energy $E_a = 0.45$ eV, a room temperature extrapolated electron mobility of $\mu = 2.8 \times 10^{-7} \text{ cm}^2/\text{Vs}$ was calculated, ¹³ leading to an electron concentration of $n = 5.5 \times 10^{17}$ cm⁻³. A low electron mobility, resulting form small polaron formation, 13 accompanied by the high charge carrier concentration in the CeO₂ thin film should certainly effect its film properties as a gate dielectric, due especially to the ease with which space charges may be activated inside the oxide. Since this space charge has a specific relaxation mechanism, its polarization response should be strongly frequency dependent, as is the case with these Pt/Si(100)/CeO₂/Pt structures. Indeed, there was a clear local maximum in loss angle tan $\delta = 3.8$ at the frequency of 1 kHz and gate bias of around 0.6 V, whereas at 100 Hz, the loss angle was already indefinite, and at 100 kHz, the loss angle increased gently around a value of 0.6 without any resonance maximum with increasing gate bias. It is suggested here, that in the vicinity of this kind of resonance frequency, there is a voltage dependent depletion layer formed in the CeO₂ film near the silicon substrate interface leading to the unusual form of the C-V curves presented in Fig. 3. This capacitance increases strongly with decreasing DC bias under 0.6 V, which becomes evident also in C-V curves measured at 1 kHz in Fig. 3, behavior similar to that of the space-charge capacitance in back-electrode Schottky contact.

4. Conclusions

The fabrication and characterization of the Pt/Si/ CeO_2/Pt MOS capacitor structure were described. The electrical properties of the CeO_2 thin films were characterized by examining and modeling the C-V response of the MOS structure utilizing conventional impedance spectroscopy. The CeO_2 layer was found to have a dielectric constant of $\varepsilon_r = 23$. A clear shift of the flat band voltage $V_{\rm FB}$ to the higher gate bias voltages with increasing frequency reflected a clear indication of the

contribution of interface charge traps. In addition to the conventional C–V measurement response, there were also some unusual features observed at low frequencies. Above a certain positive gate bias, there is a strong increase in the value of capacitance $C_{\rm ox}$ in accumulation, which is most likely due to the space charge generation and depletion at higher DC bias voltages in the CeO₂ thin film. This is consistent with our findings of high electron densities in the CeO₂ films.

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