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Size effects in Fe³⁺-doped PbTiO₃ nanocrystals—Formation and orientation of $(Fe'_{Ti}-V_O^{\bullet\bullet})^{\bullet}$ defect-dipoles

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Abstract

A set of Fe³⁺-modified PbTiO₃ nanopowders has been synthesized according to the combined polymerization and pyrolysis (CPP) route of metallorganic precursors [Erdem E, Böttcher R, Semmelhack H-C, Gläsel H-J, Hartmann E, Hirsch D. Preparation of lead titanate ultrafine powders from combined polymerisation and pyrolysis route. *J. Mater. Sci.* 2003;**38**:3211–7] with subsequent calcination at various temperatures. X-ray diffraction verifies the formation of a PbTiO₃ perovskitic phase and Raman-spectroscopy proves the existence of ferroelectricity.

Furthermore, the prevailing defect structure has been investigated by means of electron paramagnetic resonance (EPR) spectroscopy. The EPR results clearly indicate marked size effects by approaching to the critical grain size ($d_{crit} < 12 \text{ nm}$) at which a size-driven tetragonal-to-cubic phase transition is observed at room temperature. As a function of mean grain size, either ($Fe'_{Ti} - V_0^{\bullet \bullet}$) $^{\bullet}$ defect dipoles or 'isolated' defects (Fe'_{Ti} , $V_0^{\bullet \bullet}$) are formed. These results are analyzed in terms of a *core-shell model*. Accordingly, the obtained Fe^{3+} -modified PbTiO₃ nanoparticles consist of a ferroelectric core, a distorted interface region, and a cubic dead layer which is paraelectric. © 2009 Elsevier Ltd. All rights reserved.

Keywords: Size effects; Ferroelectric-to-paraelectric phase transition; Critical size; Lead titanate ferroelectrics; Iron doping; Oxygen vacancies; Defect dipoles

1. Introduction

Owing to its intrinsic high tetragonality and high Curie temperature T_C , lead titanate (PbTiO₃, PT) is a promising candidate for nano-technological ferroelectric devices. More general, ferroelectricity is a collective phenomenon involving the cooperation of polar distortions through both, short-range chemical and long-range dipolar interactions. As a consequence, the ferroelectric transition is strongly affected or even suppressed for small particle sizes (nano-structures) or confined geometries (thin films). As a function of reduced size, the long-range forces become weaker, and *vice-versa* the short-range forces become dominand. Consequently, below a critical size $d_{\rm crit}$ a size-driven para-to-ferroelectric phase transition is observed. However, the exact size of $d_{\rm crit}$ for different systems is still controversially discussed. 2,3,4,5 Typically, for thin-films critical sizes well below 10 nm are reported. 6,7,8

A still open question, however, is how size or confined geometry impacts the defect structure, such as the formation or orientation of acceptor-oxygen vacancy defect dipoles.^{9,10} In 'bulk' ferroelectrics, it was recently reported that acceptor-type centers form defect dipoles with charge-compensating oxygen vacancies, such as $(Fe'_{Zr,Ti}-V_O^{\bullet \bullet})^{\bullet}$ 11,12 or $(Cu''_{Zr,Ti}-V_O^{\bullet \bullet})^{\times}$ 13,14 in Pb[Zr_xTi_{1-x}]O₃ (PZT). By doping over the solubility limit, which typically is in the order of a few percent, secondary magnetic phases may be formed. 15 On the other hand, donor-type centers tend to form rather 'isolated' functional centers where the excess charge in form of lead vacancies (V''_{Ph}) is compensated in distant coordination spheres. 16 Concerning the $(Fe'_{ZrTi}-V_O^{\bullet \bullet})^{\bullet}$ defect dipoles, their orientation with respect to the one of the spontaneous polarization is a function of crystal symmetry and lattice distortion. 19,17,18 For this reason a marked effect on the formation and orientation of $(Fe'_{Ti}-V_O^{\bullet\bullet})^{\bullet}$ defect dipoles is expected for PbTiO₃ nanocrystals. Furthermore, upon the application of a suitable electric field the defect dipoles reorient along the direction of spontaneous polarization.²⁰

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On the other hand, it is known that the Curie temperature for 'bulk' ferroelectric systems may be altered by doping with different transition-metal ions. ²¹ Consequently, the critical size may also be impacted. For Cr^{3+} -modified PbTiO₃ nanocrystals, a critical size of $d_{\rm crit}=6$ nm has recently been reported. ²² In this work, we therefore aim to synthesize and systematically study of the interplay between grain size and characteristics of $(Fe'_{Zr,Ti}-V_0^{\bullet\bullet})^{\bullet}$ defect dipoles in Fe^{3+} -modified PbTiO₃ nanocrystals.

2. Experimental

2.1. Synthesis

The Fe³⁺-modified PbTiO₃ nanocrystals were prepared by the recently introduced combined polymerization and pyrolysis (CPP) route of metallorganic precursors, ²³ which can be applied for preparing doped perovskitic ultra-fine powders. 24,25,22 The main advantages of CPP route are the easy incorporation of transition-metal ions, the requirement of comparatively low crystallization temperatures (below 400 °C) and the simplicity of controlling the particle size by adjusting the preparation temperature. This method gives good surface morphology, phase purity and non-agglomerated nanoparticles.²³ The drawback of this route is, however, that no PbTiO₃ nanoparticles with mean grain size below 30 nm can be synthesized. In order to circumvent this shortcoming, short-time soft milling of a parent nanopowder obtained from a tempering route at medium temperature has been subsequently performed, giving particles with mean grain size below 10 nm. With the very low mechanical energy requirement it is possible to reduce the particle sizes below 10 nm.²⁷

In detail, the CPP route is based on a route for synthesizing of nanoferrites 26 and involves the following two stages: (i) preparation of metallo-organic precursor and (ii) polymerisation and pyrolysis. The chemicals used in these processes are lead(II) oxide (PbO)—Aldrich, titanium(IV) isopropoxide (Ti(O-C₃H₇)₄)—Aldrich, and methacrylic acid (CH₂=C(CH₃)-COOH)—Aldrich. Due to the synthesis conditions (open system), a certain loss of the lead component had to be taken into account, and for a corresponding compensation 10 percent surplus of PbO was used with respect to the stoichiometry. Iron(III) acetylacetonate (Fe(CH₃COCH=C(O $^{\bullet}$)CH₃)₃)—Aldrich, was employed as dopant because it does not only contain Fe³⁺ ions but is also readily soluble in chloroform.

(I) Metallo-organic precursor: Lead titanyl methacrylate was obtained as a monomeric metallo-organic precursor from lead oxide, titanium(IV) isopropoxide and methacrylic acid in boiling pure chloroform (CHCl₃) (magnetic stirring, 3 h, $150\,^{\circ}$ C) by four reactions as described below. A stabiliser (4-methoxyphenol) was added to the solution to avoid polymerisation of methacrylic acid during the heating. Individual formation reactions of lead methacrylate and titanium propoxide methacrylate from PbO and Ti(O-C₃H₇)₄can be written as

follows:

$$PbO + 2CH2=C(CH3)-COOH$$

$$\rightarrow Pb(OOC-C(CH3=CH2))2 + H2O.$$
 (1)

$$\begin{split} & \text{Ti}(\text{O-C}_3\text{H}_7)_4 + 2\text{CH}_2 \text{=-C}(\text{CH}_3 \text{--COOH} \\ & \rightarrow \text{Ti}(\text{O-C}_3\text{H}_7)_2)((\text{OOC-C}(\text{CH}_3) \text{=-CH}_2)_2) + 2\text{C}_3\text{H}_7\text{OH}. \end{split} \tag{2}$$

In presence of both PbO and Ti $(O-C_3H_7)_4$ in methacrylic acid, the reaction product will be a mixture of (1) and (2), thus:

PbO + Ti(O-C₃H₇)₄ + 4CH₂=C(CH₃)-COOH

$$\rightarrow$$
 PbTi(O-C₃H₇)₂(OOC-C(CH₃)=CH₂)₄
+ H₂O + 2C₃H₇OH. (3)

The final organometallic precursor (4) can be obtained by addition of water while refluxing. This step is necessary to complete the partial ligand exchange reaction of titanium in (2).

PbTi(O-C₃H₇)₂(OOC-C(CH₃)=CH₂)₄ + H₂O

$$\rightarrow$$
 PbTiO(OO-C(CH₃)=CH₂)₄ + 2C₃H₇OH. (4)

The solvent used in (1)–(4) was chloroform (boiling point: $62\,^{\circ}$ C) which was removed by rotary evaporation. The solid precursor was placed in an oven for 2 days at $70\,^{\circ}$ C for drying and grinded to fine powder for further use.

(II) Combined polymerisation and pyrolysis: This process is rather complex and may be described by three reaction steps:

$$n$$
PbTiO(OOC-C(CH₃)=CH₂)₄
 \rightarrow (PbTiO(OOC-C(CH₃)-CH₂)₄) _{n} (5)

$$PbCO_3 + TiO_2 \rightarrow PbTiO_3 + CO_3. \tag{7}$$

The solid state polymerisation (5) and the concomitant pyrolysis (6) proceed simultaneously at temperatures above 165 °C. Nucleation (6) and growth of PbTiO₃ particles (7) start at about 385 °C, where the temperature and the reaction atmosphere determine the resulting particle size. In Fig. 1 the corresponding calcination scheme for the control of the resulting mean particle size of the PbTiO₃ nanaocrystals. Typically, the samples were annealed for 6 h at the target temperature and hereby an increase of the particles size was observed when the annealing temperature was elevated.

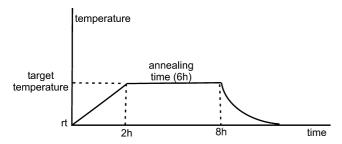


Fig. 1. Schematic representation of the calcination scheme for the CPP route in order to control the resulting particle size.

2.2. Characterization

X-ray diffraction (XRD) measurements were carried out at room temperature with a powder diffractometer Huber G670 equipped with a germanium monochromator, using Cu- K_{α} radiation. The particle size was determined using Scherrer's equation $D = (k\lambda/\beta\cos(\theta))$ where D is the mean grain size, k the Scherrer's constant (k = 0.9 for FWHM), λ the X-ray wavelength, β the FWHM of a diffraction peak, and θ the diffraction angle.

The FT-Raman spectra were recorded using an Bruker RFS 100 spectrometer with Nd:YAG laser as excitation source at a power level of up to 30 mW.

The X-band (9.5 GHz) continuous-wave (cw) EPR measurements were performed using an ESP 300E spectrometer (Bruker), equipped with a rectangular TE₁₁₂ resonator. All EPR measurements performed at ambient temperature.

For numerical spectrum analysis of the EPR data, the free trivalent iron ion is considered as possessing five unpaired electrons in the 3d shell $(3d^5)$ with a high-spin $^6S_{5/2}$ ground state (S=5/2). The crystal field originating from the nearest neighbor oxygen ions (and possibly oxygen vacancies) lifts the Fe³⁺ spin degeneracy resulting in three Kramer's doublets, whose degeneracy in turn is lifted by an external magnetic field. Neglecting hyperfine interaction, as the only isotope with non-vanishing nuclear spin 57 Fe is present in only 2.15% natural abundance, an approximate spin Hamiltonian for these high-spin system can be written as

$$\mathcal{H} = g_{iso}\beta_e \mathbf{B}_0 \cdot \mathbf{S} + \sum_{k}^{-k \le q \le k} B_k^q O_k^q(S_x, S_y, S_z)$$
 (8)

where the electronic g-matrix may be replaced by an isotropic value of $g_{iso}^{{\rm Fe}^{3+}}=2.002$, β_e denotes the Bohr magneton, ${\bf B}_0$ the external magnetic field, B_k^q are the FS spin-Hamiltonian parameters, and O_k^q are the extended Stevens spin operators. In principle, all conceivable B_k^q parameters, resulting from 2nd to 4th-(Fe³⁺) tensors, have to be considered. However, experimentally only the 2nd-rank B_2^q parameters for Fe³⁺ could reliably be determined by means of numerical spectrum simulation 30 using Eq. (8).

3. Results and discussion

In order to verify the formation of the expected tetragonal $PbTiO_3$ perovskite phase after calcination, XRD data of the CPP

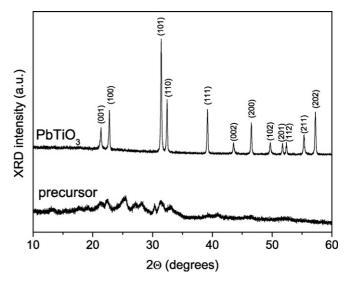


Fig. 2. Powder diagram of a PbTiO $_3$ sample prepared by the combined polymerization and pyrolysis (CPP) route and calcined at 800 $^\circ$ C (top) in comparison with that of the precursor (bottom). The indexed reflections indicate the formation of single-phase tetragonal PbTiO $_3$.

precursor and the resulting PbTiO₃ nanopowder after calcination at 800 °C are shown in Fig. 2. Whereas the XRD pattern of the precursor shows no significant peaks, indicating the expected amorphous structure, the XRD pattern of a calcined sample exhibits the pattern expected for tetragonal PbTiO₃powder with the corresponding Bragg reflections indexed.

By subsequently using Raman spectroscopy it is possible to account for the expected soft-mode behavior of ferroelectric PbTiO₃ particles. The soft mode E(1TO) is a prerequisite for the existence of long-range phenomena and often is used as a fingerprint for the existence of ferroelectricity. The corresponding spectra are displayed in Fig. 3. For a sample calcined at 800 °C, all expected Raman-active modes for PbTiO₃²⁹ were observed. As compared to the Raman frequencies reported for

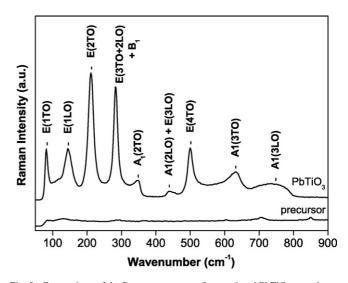


Fig. 3. Comparison of the Raman spectrum of a powdered PbTiO $_3$ sample synthesized via combined polymerization and pyrolysis (CPP) route and calcined at 800 °C (top) with that of the precursor (bottom). All Raman modes expected for PbTiO $_3$ could be observed and assigned.

'pure' PbTiO₃ no significant shifts could be registered for the Fe³⁺-modified PbTiO₃ compounds. On the other hand, for the amorphous precursor no significant Raman-active modes could be detected.

The 'method-of-choice' for the characterization of prevailing defect structure in the nanoparticles with respect to the formation of $(Fe'_{Ti}-V_O^{\bullet \bullet})^{\bullet}$ defect dipoles, as well as its orientation as a function of PbTiO3 mean grain size is electron paramagnetic resonance (EPR) spectroscopy. 9,10 The information content obtained from the analysis of the EPR spectra is twofold. First, it can be distinguished between 'isolated' Fe'_{Ti} functional centers and centers that have formed a $(Fe'_{Ti}-V_O^{\bullet \bullet})^{\bullet}$ defect dipole for reasons of charge compensation. Second, the site symmetry at the Fe³⁺-site is obtained, for which reason in case of a $(Fe'_{Ti}-V_O^{\bullet \bullet})^{\bullet}$ defect dipole, its orientation with respect to the orientation of the spontaneous polarization may be deduced. 19 Correspondingly, the defect dipoles may be oriented either parallel $(Fe'_{Ti} - V_O^{\bullet \bullet})^{\bullet}_{||}$ or perpendicular $(\mathrm{Fe}'_{\mathrm{Ti}} - V_{\mathrm{O}}^{\bullet \bullet})^{\bullet}_{\perp}$ with respect to the orientation of P_S . In case of an 'isolated' Fe'_{Ti} center, an isotropic EPR line accounts for a cubic and paraelectric PbTiO₃ phase.

The corresponding EPR spectra for the obtained Fe³+modified PbTiO₃ nanopowders with varying mean grain size are given in Fig. 4. For comparison, also the EPR spectrum of a 'bulk' sample with mean grain size of 500 nm is depicted (top). The spectrum consists of a dominant low-field peak at $\approx 100\,\mathrm{mT}$ characteristic for the existence of $(\mathrm{Fe'_{Ti}}-V_\mathrm{O}^{\bullet\bullet})^{\bullet}{}_{||}$ defect dipoles. Por reduced grain size, additional peaks occur in the EPR spectra (center, bottom). These may be explained by the following scenario. For nano-sized samples an increasing contribution arises from the surface of the nanograins. According to a *core*–shell model, 32,22 this surface has cubic PbTiO₃ structure and thus is paraelectric. In cubic crystal symmetry, the advantage in energy of formation of an $(\mathrm{Fe'_{Ti}}-V_\mathrm{O}^{\bullet\bullet})^{\bullet}$ defect associate over the 'isolated' defects $(\mathrm{Fe'_{Ti}},\ V_\mathrm{O}^{\bullet\bullet})^{11}$ vanishes. Consequently,

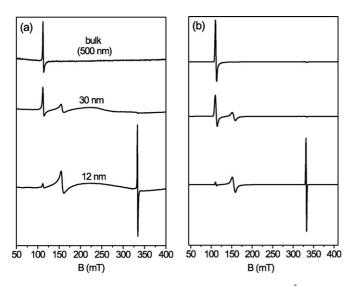


Fig. 4. Room-temperature X-band (9.5 GHz) EPR spectra of Fe^{3+} -modified PbTiO₃ nanopowders with varying mean grain size. (a) Experiment and (b) numerical simulation.

the observation of an isotropic EPR at g=2.002 (Fig. 4, bottom) accounts for the existence of an isolated Fe'_{Ti} center in a paraelectric surface region. The third resonance at $\approx 150\,\mathrm{mT}$ is characteristic for a $(\mathrm{Fe'_{Ti}}-V_0^{\bullet\bullet})^{\bullet}$ defect dipole of rhombic site symmetry at the site of the Fe³⁺ functional center. Its origin may be twofold. First, it may be a $(\mathrm{Fe'_{Ti}}-V_0^{\bullet\bullet})^{\bullet}_{\perp}$ defect dipole. Second, it may be a $(\mathrm{Fe'_{Ti}}-V_0^{\bullet\bullet})^{\bullet}$ defect dipole in a distorted crystal symmetry. Both scenarios agree with the existence of a distorted interface layer in the nanocrystals between the ferroelectric core and the paraelectric surface.

However, the present results do not allow to monitor the complete size-driven phase transition, for which exclusively 'cubic' EPR spectra are expected. Consequently, only an upper limit for the corresponding critical size $d_{\rm crit} < 12$ nm can be given.

The obtained EPR spectra may nicely be reproduced by means of numerical spectrum simulation 30 by assuming a superposition of three different centers with axial, rhombic and cubic site symmetries (cf. Fig. 4(b)). The broad line occurring only in the spectra of the nano-sized compounds (30 and 12 nm) between 150 and 270 mT is due to the iron oxide (Fe₂O₃) on the surface of the nanograins and was not considered in the spectrum simulations.

4. Conclusion and outlook

synthesis of Fe³⁺-modified summary, the PbTiO₃nanopowders has been reported by using the CPP route. The EPR results clearly indicate marked size effects by approaching to the critical grain size ($d_{crit} < 12 \text{ nm}$) at which a size-driven tetragonal-to-cubic phase transition is observed at room temperature. The obtained results agree well with the proposed *core-shell model* ^{32,22} and are a superposition of three phases corresponding to a ferroelectric core (tetragonal), a distorted interface layer (rhombic) and a paraelectric shell (cubic). The Fe³⁺-modified PbTiO₃ nanoparticles thus consist of a ferroelectric core, and a cubic dead layer which is paraelectric. These results are also consistent with the recent Landau phenomenological theory calculations for confined ferroelectric nanoparticles. 33

In order to enlighten the interplay between nano-structuring and defect structure in more detail, the existing set of samples has to be extended. In particular, by soft-milling compounds with mean grain sizes below the critical grain size d_{crit} shall be obtained. Furthermore, we aim to investigate the size-driven phase transition and accurately determine the critical grain size d_{crit} . Future work will be the detailed extension of structural characterisation by using mainly multi-frequency high-field EPR and the standard methods such as XRD, Raman and dielectric spectroscopy. Particularly, in the content of size effects, the soft mode reduction, change of $T_{\rm C}$ and softening of dielectric permittivity curves are the important open questions to be considered. Multi-frequency high-filed EPR (Q-band (35 GHz), W-band (95 GHz) and higher frequencies) have the advantage to detect the higher Fe³⁺ energy transitions. So that, we can have more information about the sign and the size dependency of the zero field splitting parameter.

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